

RF Electron Gun

Lecture No. 3: Operation and Applications of RF Gun

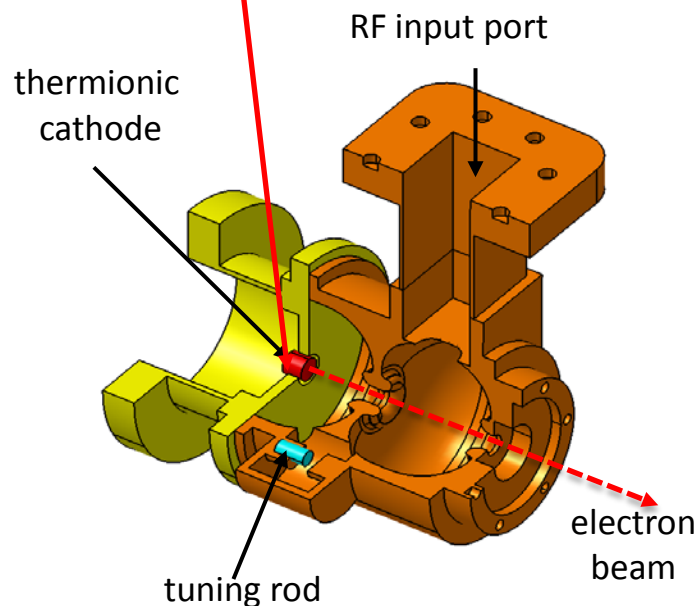
Sakhorn Rimjaem

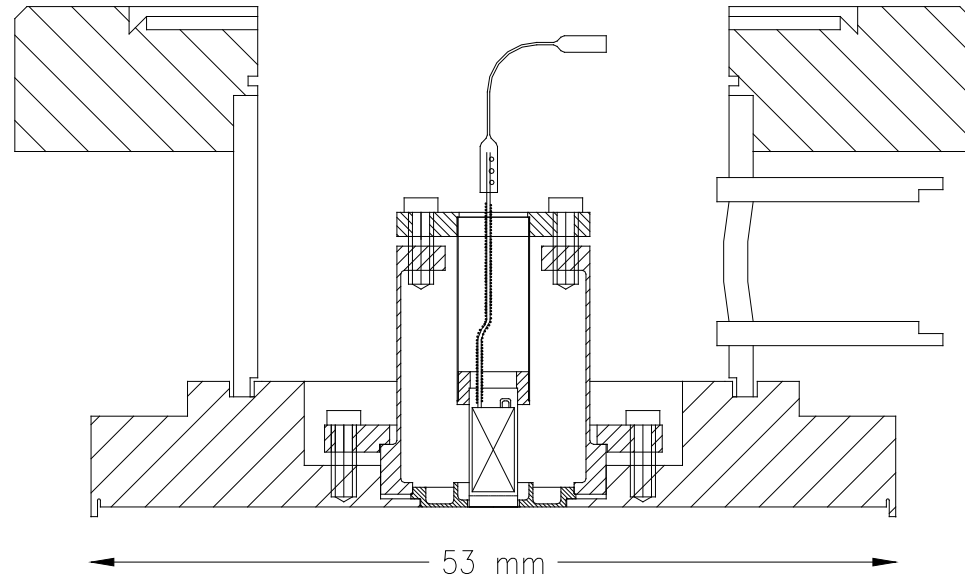
Department of Physics and Materials Science,
Faculty of Science, Chiang Mai University (CMU), Thailand

- **Operation of RF Electron Gun**
 - Cathode preparation
 - Operation conditions
 - Measurements of electron beam properties
- **Applications of Electron Sources**
 - Electron sources for high energy physics accelerator
 - Electron sources for light sources
 - Electron sources for direct applications
- **Generation and Applications of Ultra Short Electron Bunches Produced by RF Gun**
 - Generation and applications of ultra short electron bunches at Chiang Mai University

Cathode assembly

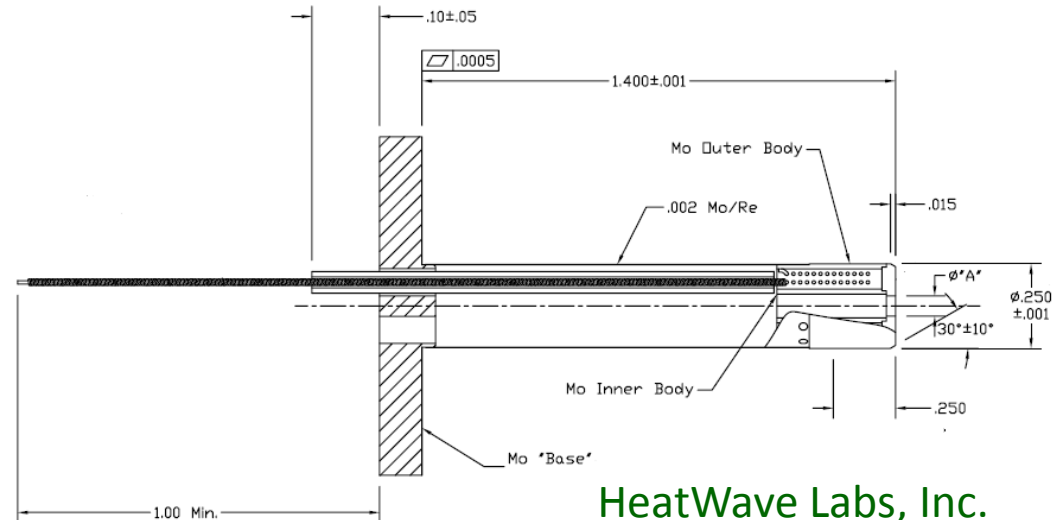
- Cylindrical SST cathode holder with screws for adjusting cathode position.
- Cathode with flat circular emitting surface of 3 mm is heated via filament connected to an external power supply.





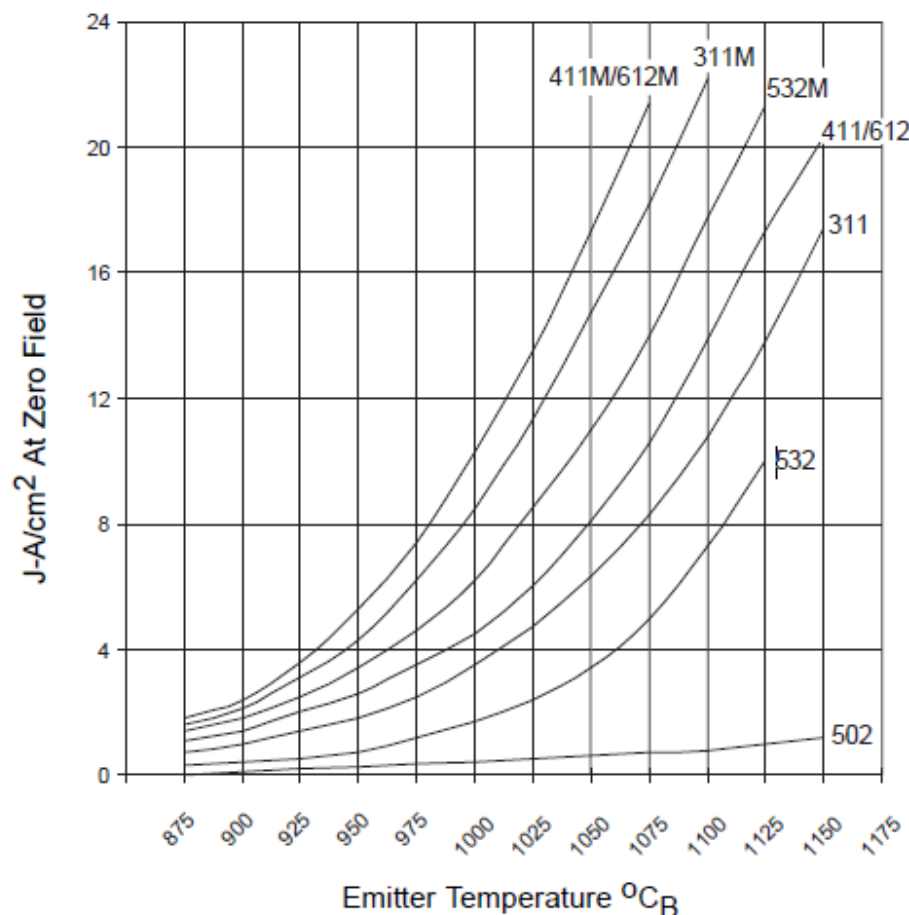
Thermionic cathode is an impregnated dispenser tungsten cathode coated with Os/Ru.

→ Flat circular emitting surface of 3 mm radius.



HeatWave Labs, Inc.

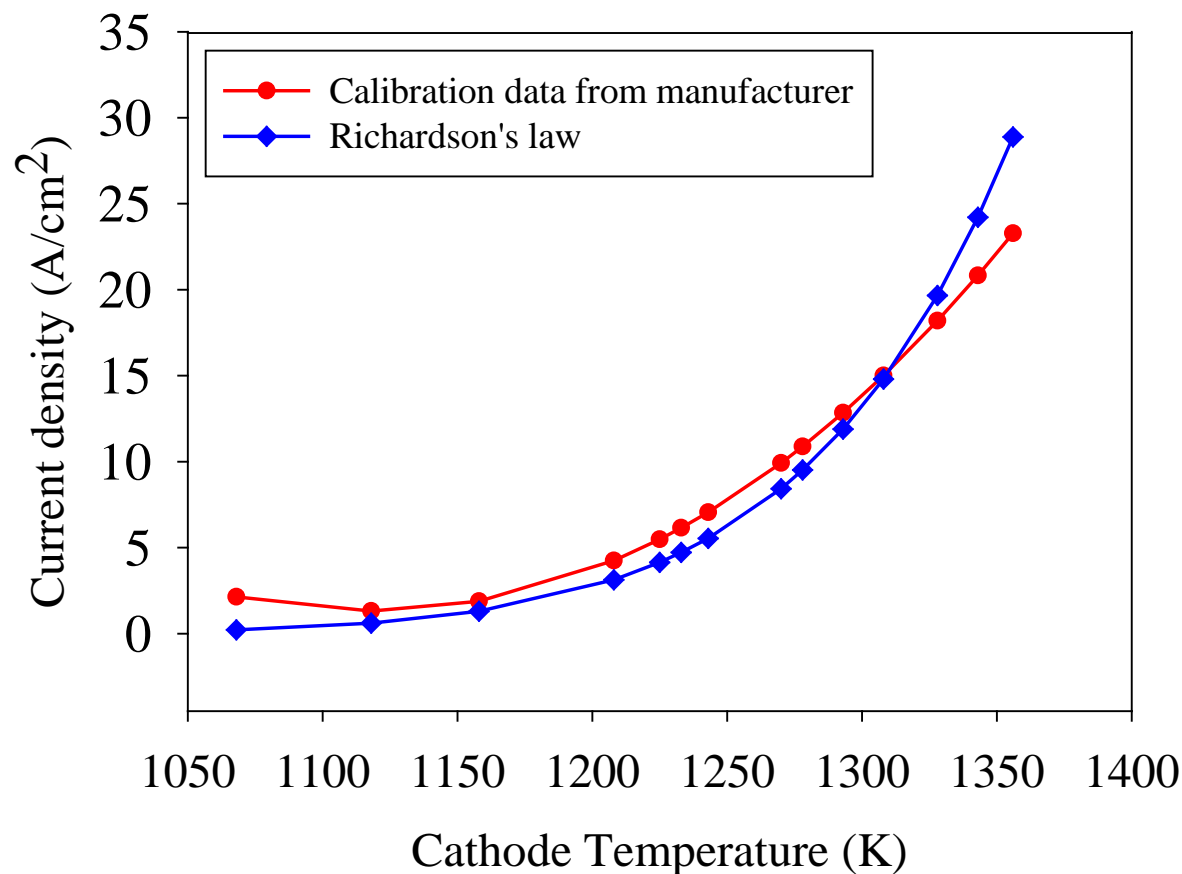
HeatWave Labs, Inc.



“Type M cathode” is an improved version of the **impregnated cathode**, which coated with an alloy of platinum group metals on the emissive surface to lower the work function.

→ It can be operated at lower temperatures than the standard dispenser cathode for the same current density, which results in longer operating lifetime.

Comparison of current density as a function of cathode temperature from manufacturer's calibration data and Richardson's law.



The cathode needs to be heated to a high temperature ($\sim 900\text{-}1000^\circ\text{C}$) for electron emission.

→ If the heat leakage is too high, the cathode might not reach the electron emission temperature.

→ Thermal radiation from hot cathode can heat up the cavity wall near the cathode to expand due to thermal expansion. This may results in a resonant frequency shift of the gun.

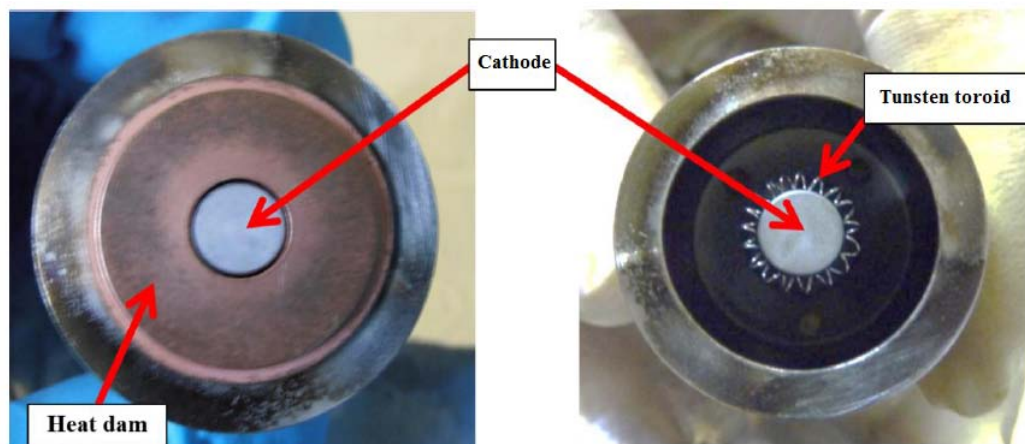
To prevent heat lost to the surrounding area, the cathode must be insulated from material.

→ A flat heat-dam made of Hastelloy C-276 with copper electroplating is used.

→ A narrow gap was made between the cathode stem and the heat-dam to function as a thermal insulator.

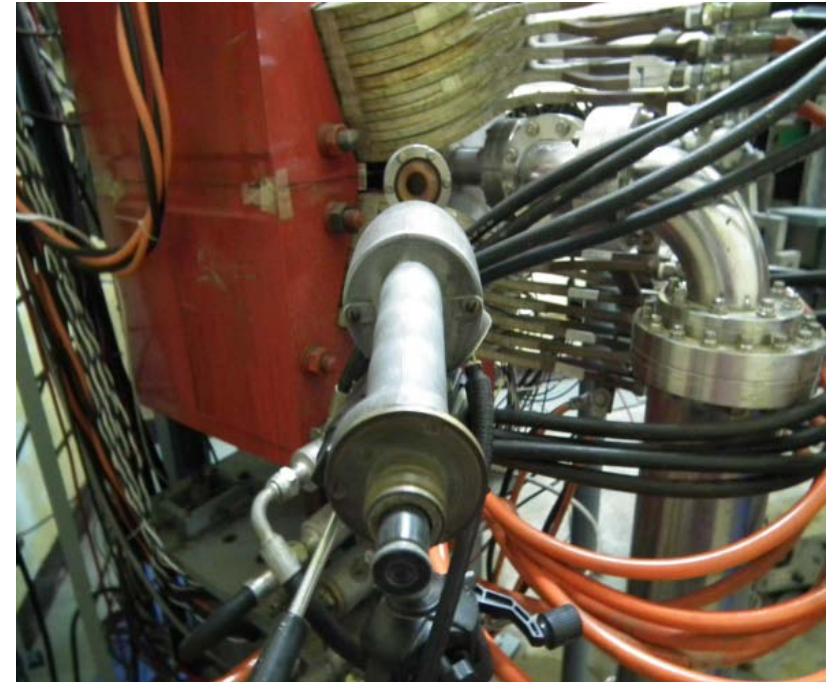
→ A toroidal- tungsten-spring is placed around the cathode stem to minimize heat leakage.

→ The **cathode temperature** depends greatly on the **number of turns of the spring**.





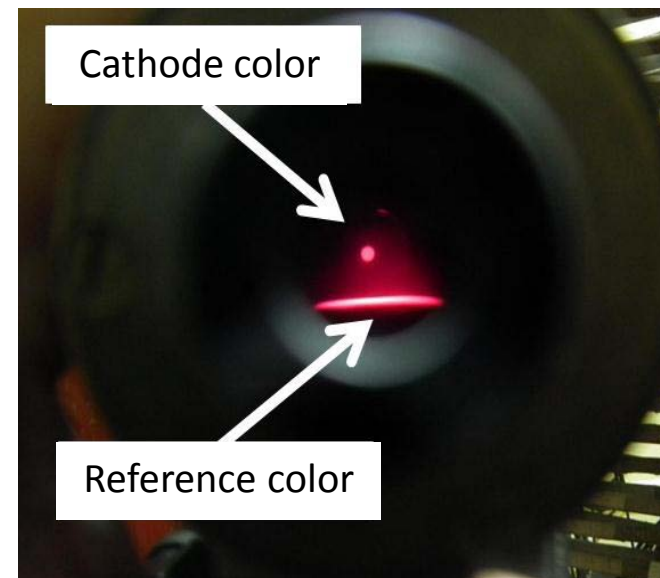
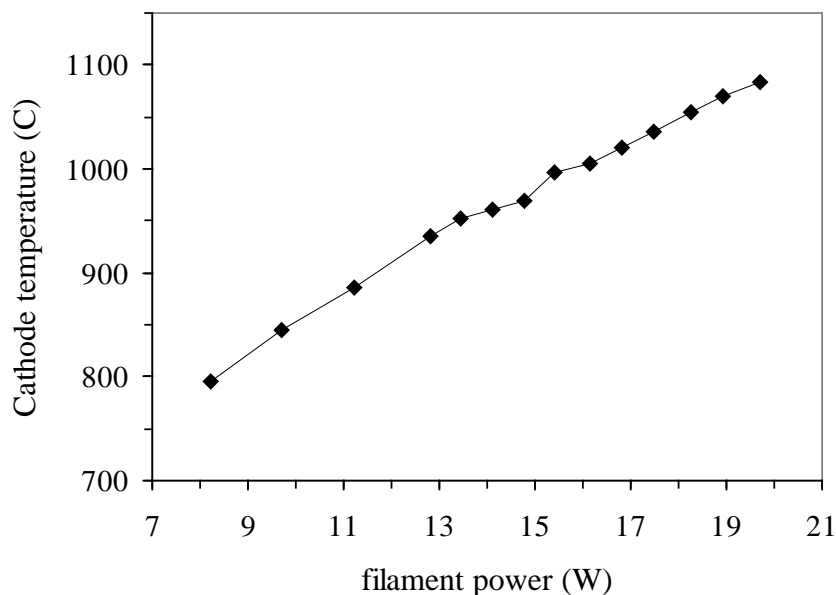
**Pyrometric measurement after
cleaning and baking**



**Pyrometric measurement for the gun,
which installed in the beam line.**

Cathode tests (Pyrometric measurements)

- To activate the cathode to temperature $> 1050\text{ }^{\circ}\text{C}$
- Measure the cathode temperature by using an optical Pyrometer
- Required for new cathode or when cathode experiences poor vacuum or chemical contamination.



**Operating temperature $\sim 900\text{-}1000^{\circ}\text{C}$
(cathode heating power $\sim 13\text{-}17\text{ W}$)**

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Input (forward) RF power given to the RF-gun is

$$P_{forward} = P_{cavity} + P_{beam} + P_{reflected}$$

Beam power is related to the kinetic energy of electron beam as

$$P_{beam} \approx (\langle KE \rangle) I_{beam}$$

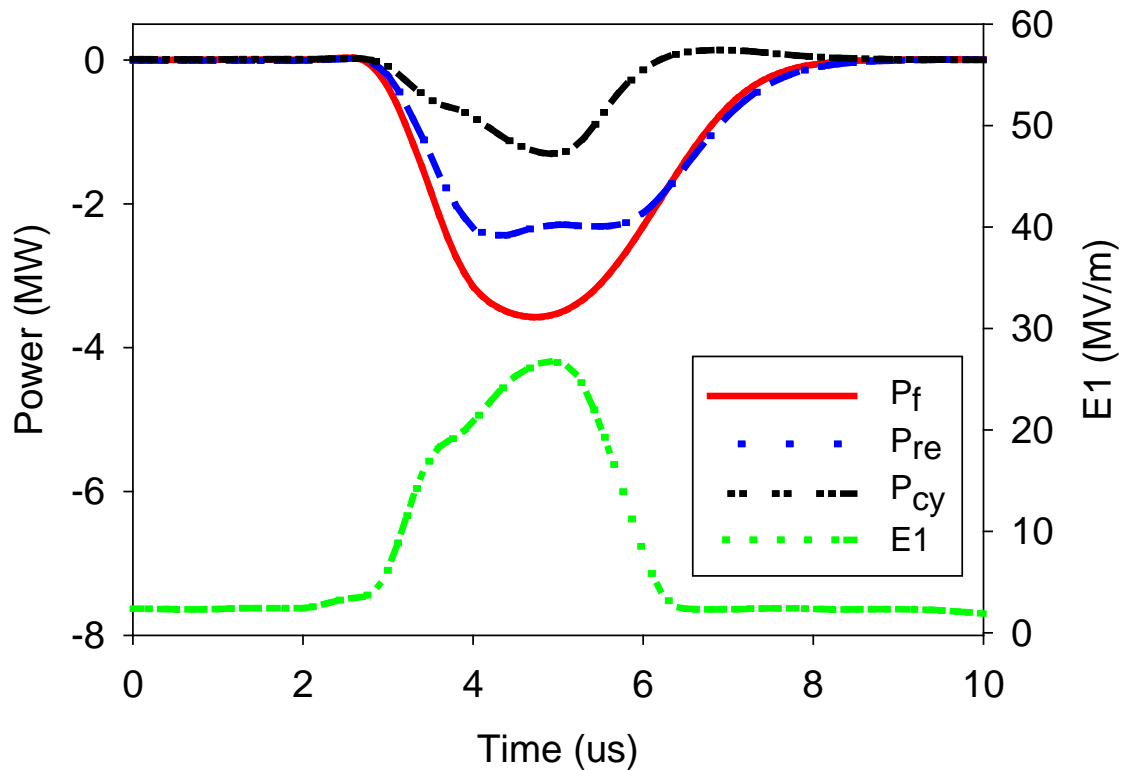
In of no electron emission from the cathode, the cavity wall losses is

$$P_{cavity} = P_f - P_{reflected}$$

Compare with theoretical cavity wall losses (which can get from simulation)

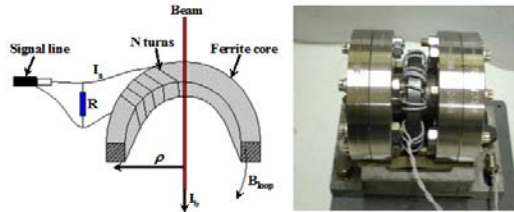
→ Then, the average electric field gradient (E_0) can be calculated:

$$P_{cavity} = \frac{V_{acc}^2}{r_s d} = \frac{E_1^2 d}{r_s} \Rightarrow E_1^2 = \frac{P_{cavity} r_s}{d}$$

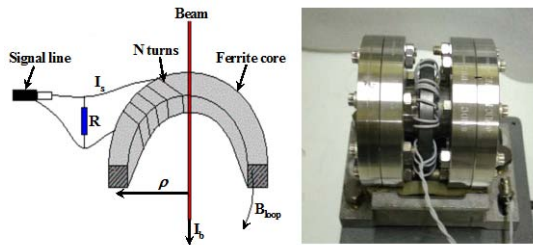


Measured RF powers and calculated average electric field gradient

Current Monitor

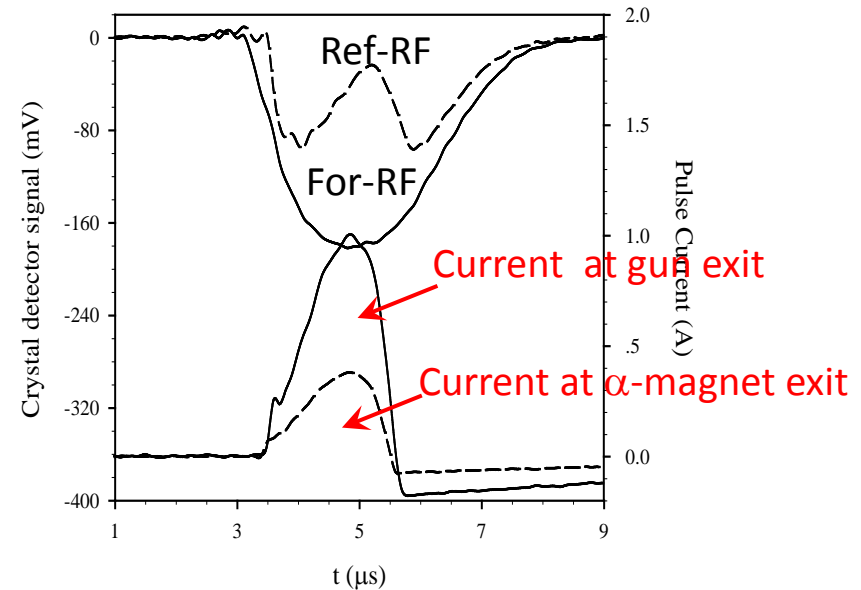


Schematic model of current transformer



Actual current transformer

$$I_b = N_s I_s = \frac{N_s}{R} V_s = \frac{8}{50} V_s = 0.16 V_s$$

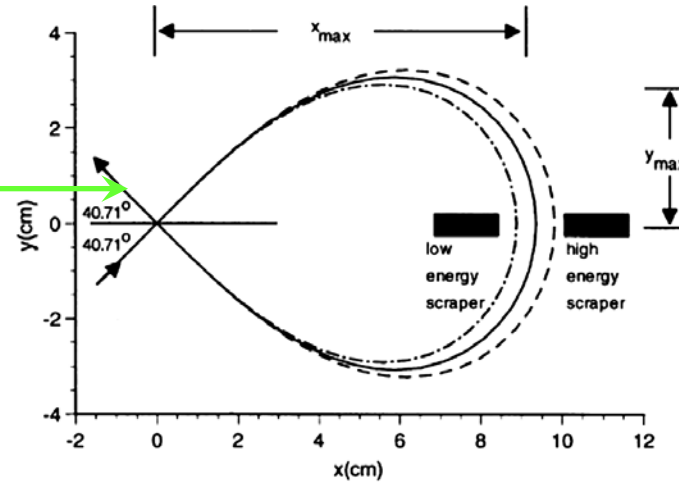


Peak current of ~ 1 A at ~2 MeV from RF-gun

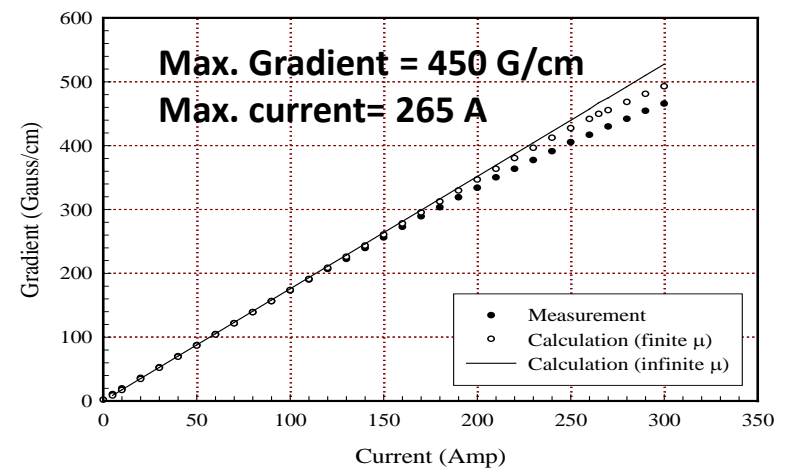
- Beam power $P_b \sim I_b \times E_{kin} = 2 \text{ MW}$
- Cavity wall losses $P_{cy} \sim 1.46 \text{ MW}$
- $P_{cy} + P_b = 3.46 \text{ MW}$

Peak current of 0.4-0.5 A at α - magnet exit
(50-60% is filtered out by the energy slits).

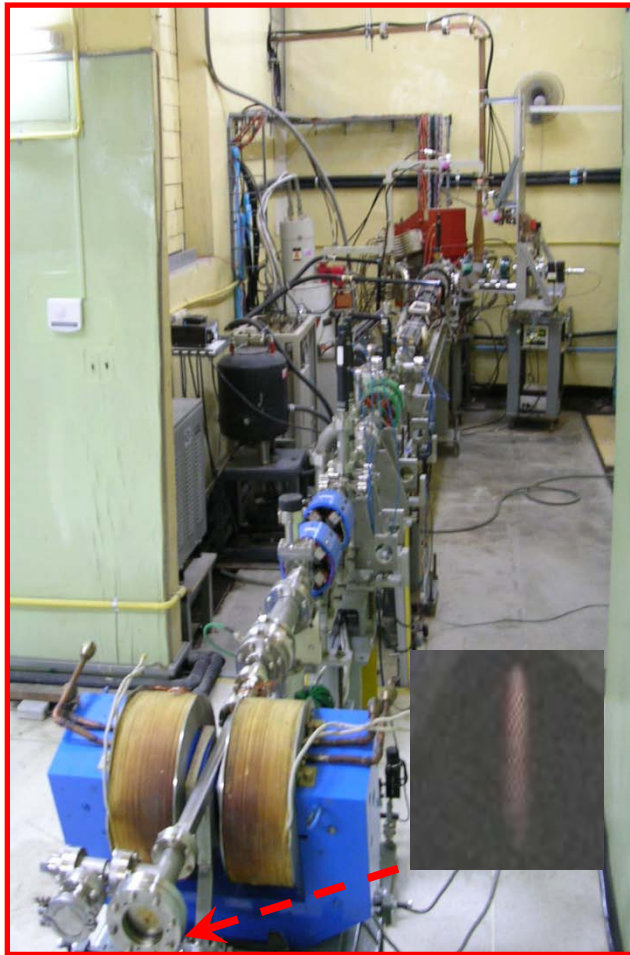
Using energy slit inside alpha magnet vacuum chamber: $E_{\max} = 2\text{-}2.4\text{ MeV}$



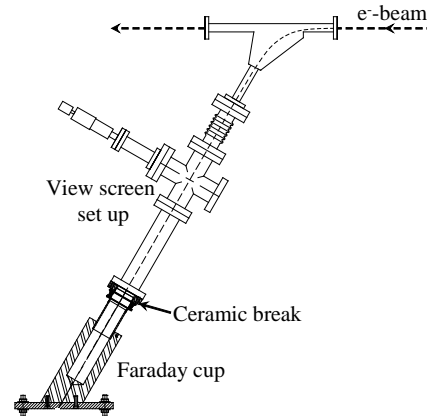
$$x_{\max} = 75.05 \sqrt{\frac{cp}{mc^2 g}}$$



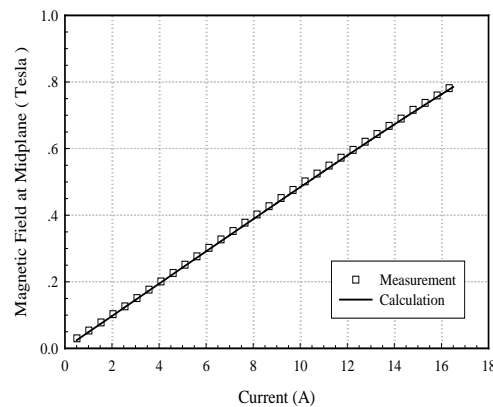
Dipole magnet is used as electron beam dump + energy spectrometer.



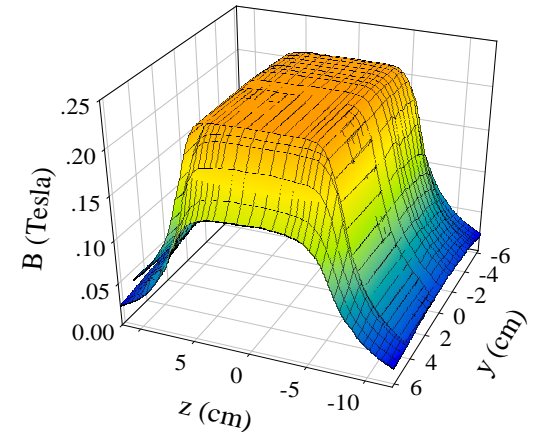
10-15 MeV electron beam



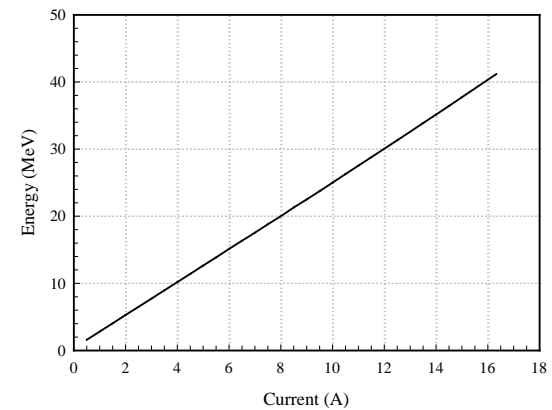
Deflect electron beam 60°
respect to beam axis



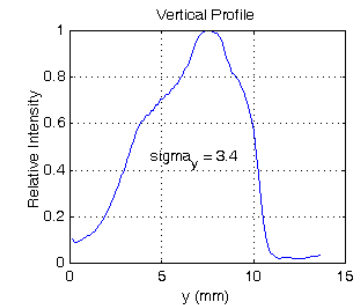
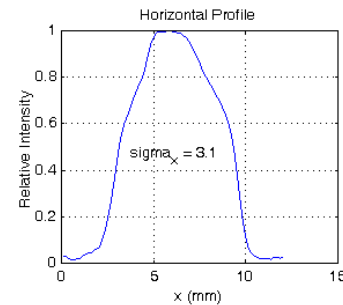
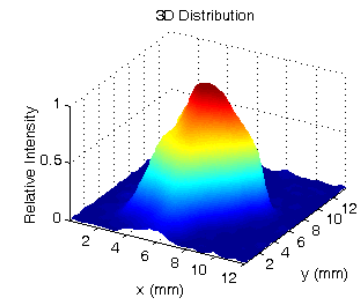
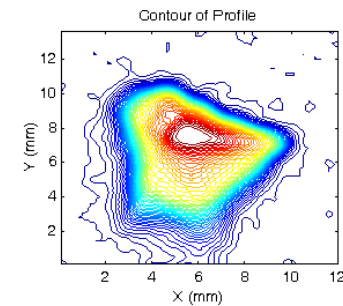
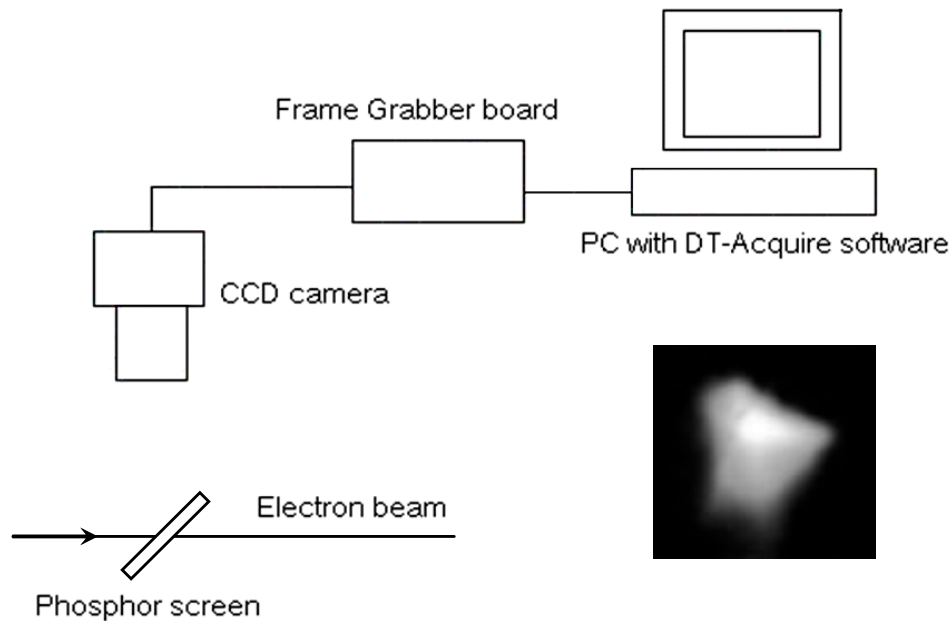
$B \sim 0.8$ Tesla at current of 16 A



Actual 3D-field distribution of
dipole magnet



$$E(\text{GeV}) = \frac{0.2998 \int B_y dz}{\beta \alpha}$$



Schematic layout of beam profile measurement setup and a 2.4 MeV electron beam image (SC2)

- Phosphor screen ($\text{Gd}_2\text{O}_3:\text{Tb}$ deposited on Al-plate)
- CCD camera
- Frame grabber board (DT3315 Data-Translation)
- PC with DT-Acquire software

Relative intensity distribution of electron beam in 2D and 3D and the horizontal and vertical beam profiles

(MATLAB code BAP, S. Chumphongphan)

High-power RF parameters, typical operating condition and electron beam-loading parameters of the PBP-CMU RF-gun.

Parameter	Value	Unit
Resonant frequency in $\pi/2$ -mode	2856	MHz
Operating temperature	27.5	°C
Input RF peak power	3.55	MW
Dissipation RF power	1.36	MW
External RF-coupling coefficient	8.3	
Maximum kinetic energy	~ 2.5	MeV
Beam current	0.7–1	A
RF-pulse length (FWHM)	~ 2.8	μs
RF repetition rate	10	Hz
Beam-pulse length (FWHM)	1–2	μs
Charge per micro-bunch	~ 0.2 -0.25	nC

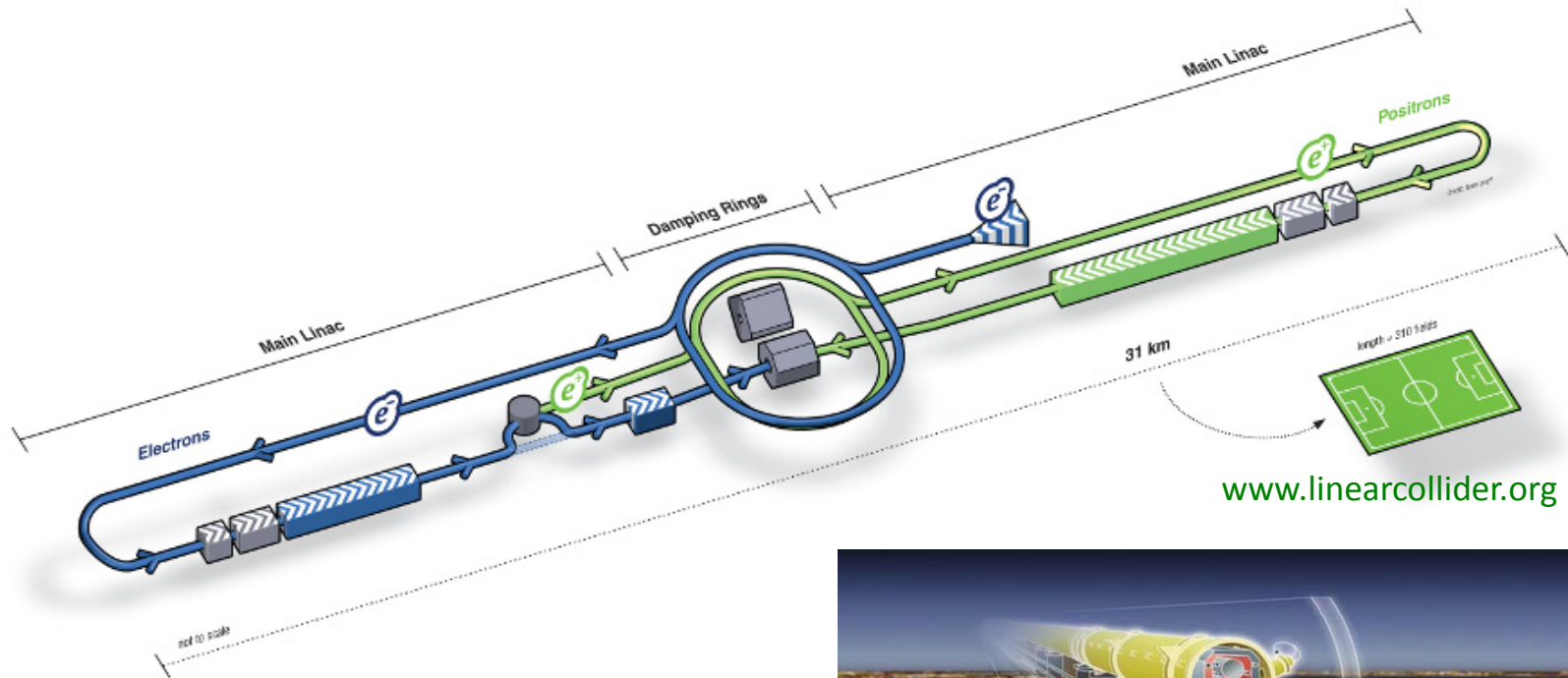
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Example: Stanford Linear Accelerator Laboratory (SLAC)

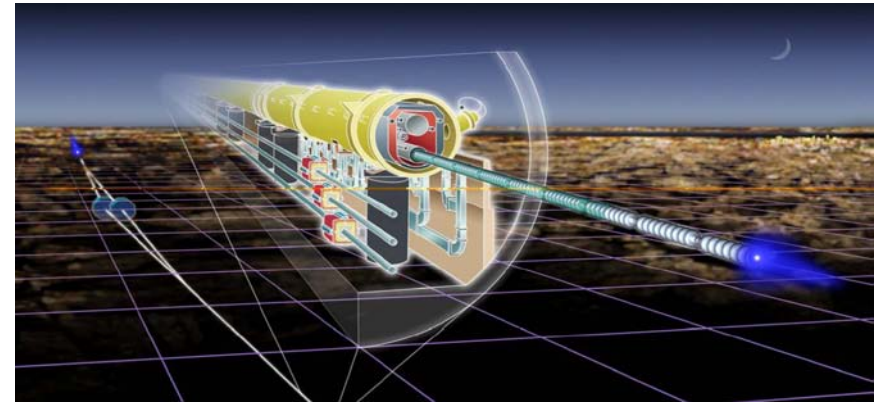


3-miles linear accelerator at SLAC can accelerate electron beam to 10 GeV.

Example: the International $e^- - e^+$ Linear Collider (ILC)

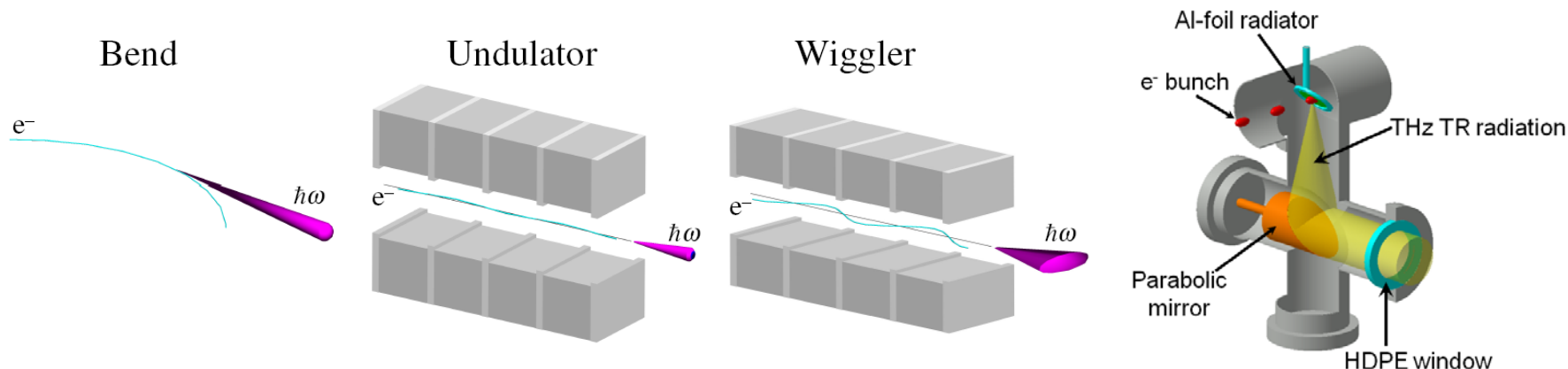


It is initially planned to have a collision energy of 500 GeV, with the possibility for a later upgrade to 1000 GeV (1 TeV).



(A. Sessler and E. Wilson, Engines of Discovery: A Century of Particle Accelerators)

Accelerating charges emit electromagnetic radiation



Radiation brightness

■ In geometric optics, spectral brightness is defined as photon flux density in phase space about a certain frequency (number of photons per unit time per unit area per unit solid angle):

$$B = \frac{d^2 I}{dA d\Omega} = \frac{d^2 I}{dx dx' dy dy'} \propto \frac{2I_p}{\varepsilon_x \varepsilon_y}$$

I - electron beam current

A - transverse area

Ω - the divergence

x, y – transverse coordinates

x', y' - indicate derivatives with respect to z

∴ We need electron beams with small emittance and high peak current.

Thermionic DC guns
Thermionic RF guns



PETRA III @ DESY (Germany)



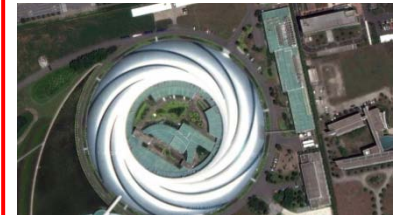
Diamond (UK)



ALBA, Spain



ESRF, France



Shanghai light source
(China)



Spring 8 (Japan)



Siam Photon (Thailand)



Australia light source
(Australia)



NSLS (USA)

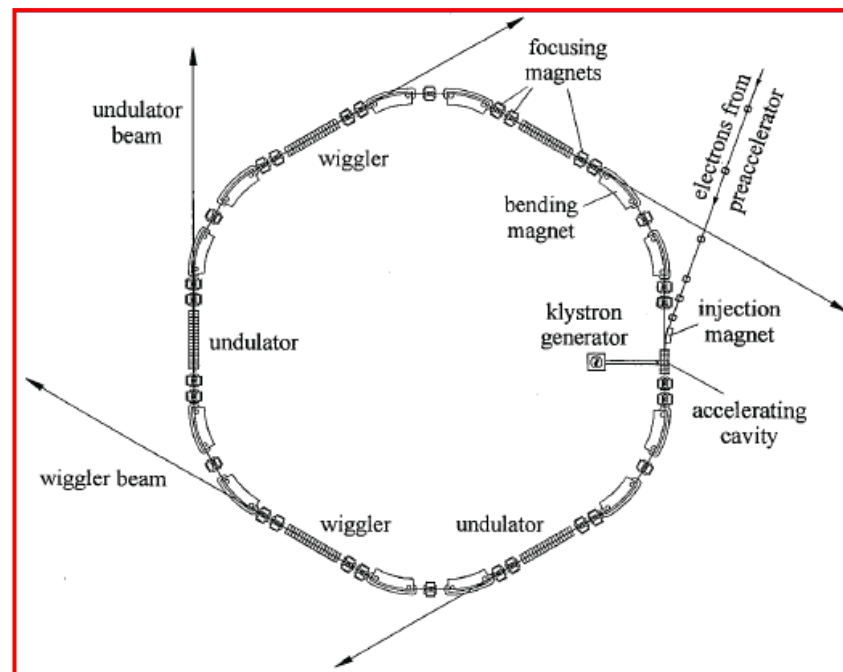
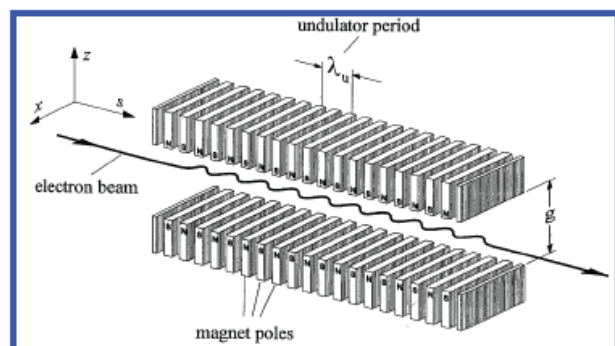
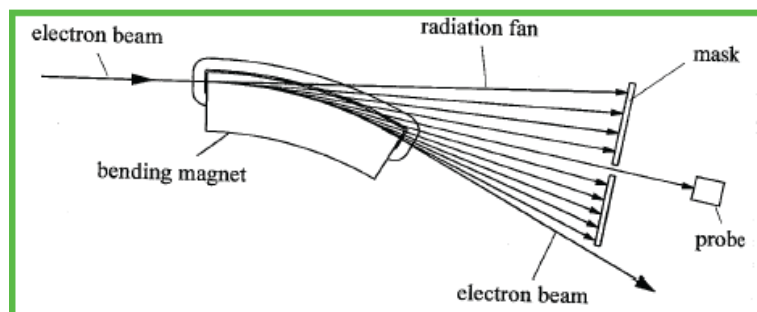


APS (USA)

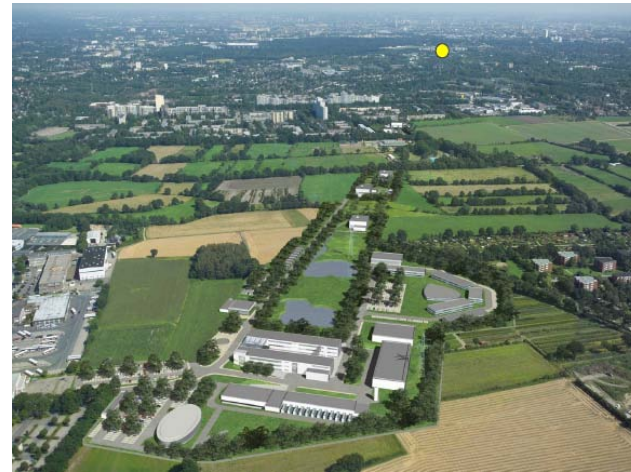
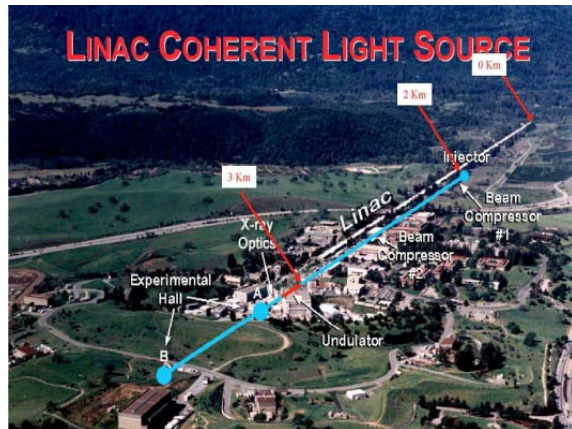


Brazilian Synchrotron

- **1st Generation** (1970s): Many HEP rings are parasitically used for X-ray production
- **2nd Generation** (1980s): Many dedicated X-ray sources (light sources)
- **3rd Generation** (1990s): Several rings with dedicated radiation devices (wigglers and undulators)
- **4th Generation** (Present): Free Electron Lasers (FELs) driven by linear accelerators



K. Wille, *The Physics of Particle Accelerators: An Introduction*, Oxford University Press, 2000.



- European XFEL, Germany ($\lambda \sim 0.1-1.6$ nm)
- FLASH @ DESY, Germany ($\lambda \sim 0.1-7$ nm)
- SPARC, Italy ($\lambda \sim 0.6-40$ nm)
- FERMI @ Elettra, Italy ($\lambda \sim 10-100$ nm)
- SwissFEL, Switzerland ($\lambda \sim 0.1-7$ nm)

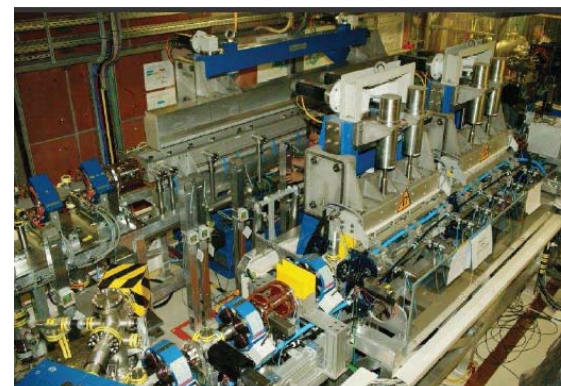
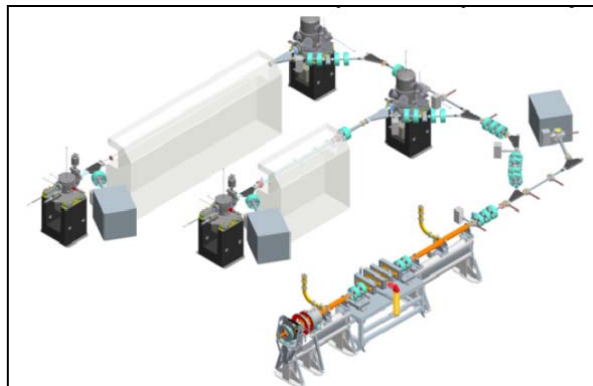


- LCLS FEL @ SLAC, USA ($\lambda \sim 0.15-1.5$ nm)
- HGHG FEL @ NSLS, BNL, USA ($\lambda \sim 193$ nm)

Photocathode RF guns
(SACLA XFEL uses DC gun)

- SACLA XFEL @ Spring-8, Japan ($\lambda > 0.1$ nm)
- Shanghai FEL
- Pohang XFEL

Fritz Haber Institute THz FEL
Max Planck Institute
Berlin, Germany



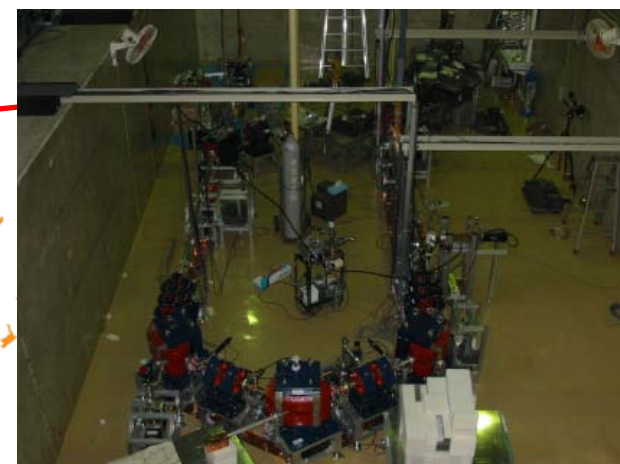
ELBE @ HZDR
Germany



UCSB FEL, Santa Barbara, USA



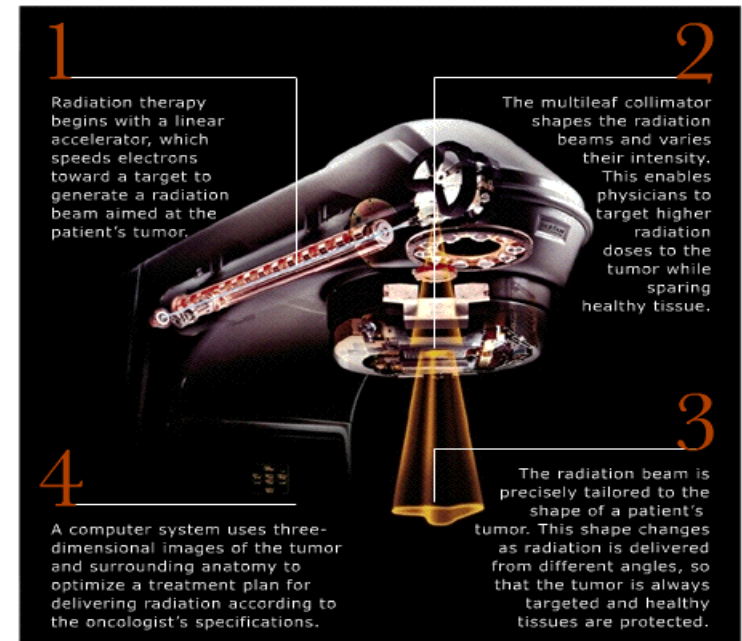
KU-FEL, Kyoto University, Japan



Thermionic DC guns
Thermionic RF guns
Photocathode RF guns



Thermionic DC guns



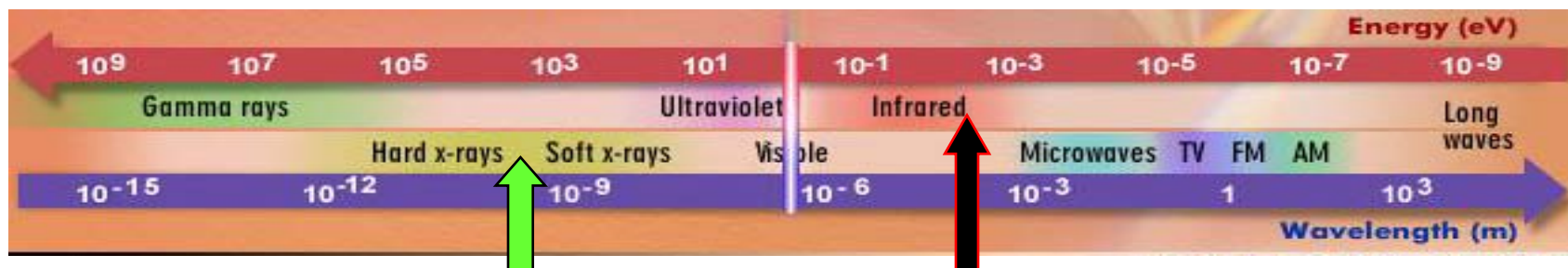
- A modern system for treating a patient with x-rays produced by a high energy electron beam.
- The whole device is mounted on a gantry. As the gantry is rotated, so is the accelerator and the resulting x-rays.
- Then the radiation can be delivered to the tumor from all directions with precise control for positioning.

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➤ Direct Applications:

- Ultrafast time-resolved electron diffraction for studying dynamic of chemical reaction (A. Zewail Noble Price Winner 2002)
- Ultrafast electron pulse radiolysis (K. Ka et al., Rev. of Sci. Instr., 2012)

➤ Production of short (femtosecond) photon pulses:



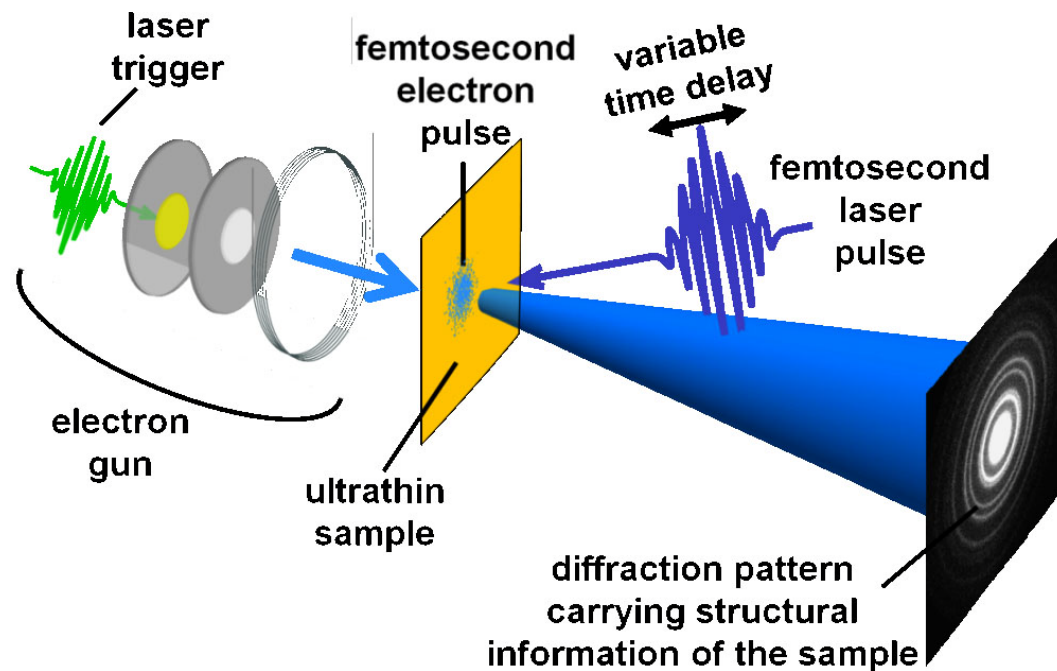
Femtosecond X-ray pulses for dynamic study at atomic scale

Intense far-infrared/THz radiation at frequencies of 100 GHz - 10 THz (λ : 3 – 0.03 mm) via, e.g.,

- Transition radiation
- Undulator radiation
- Free-electron lasers (FELs)

Examples: Photocathode RF guns

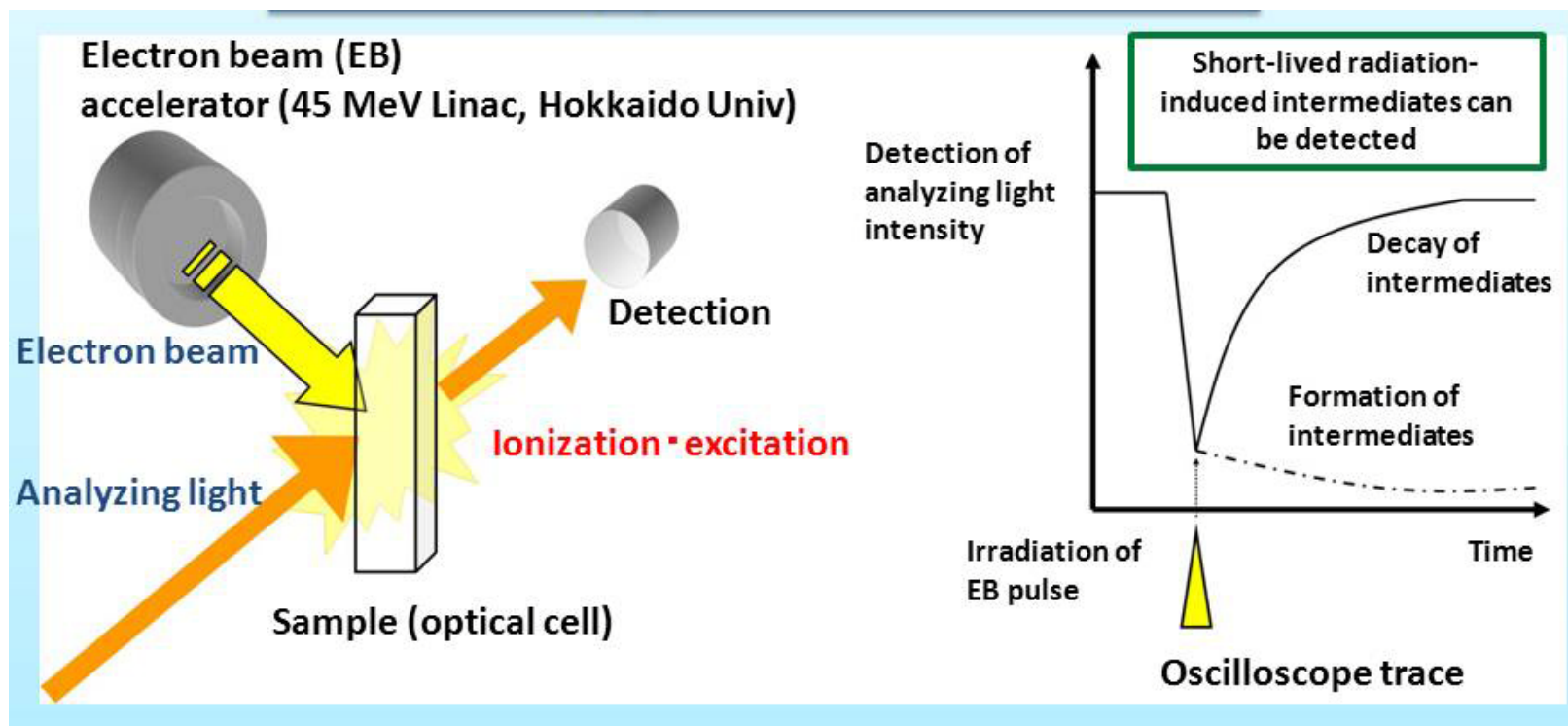
- Ultrafast time-resolved electron diffraction for studying dynamic of chemical reaction



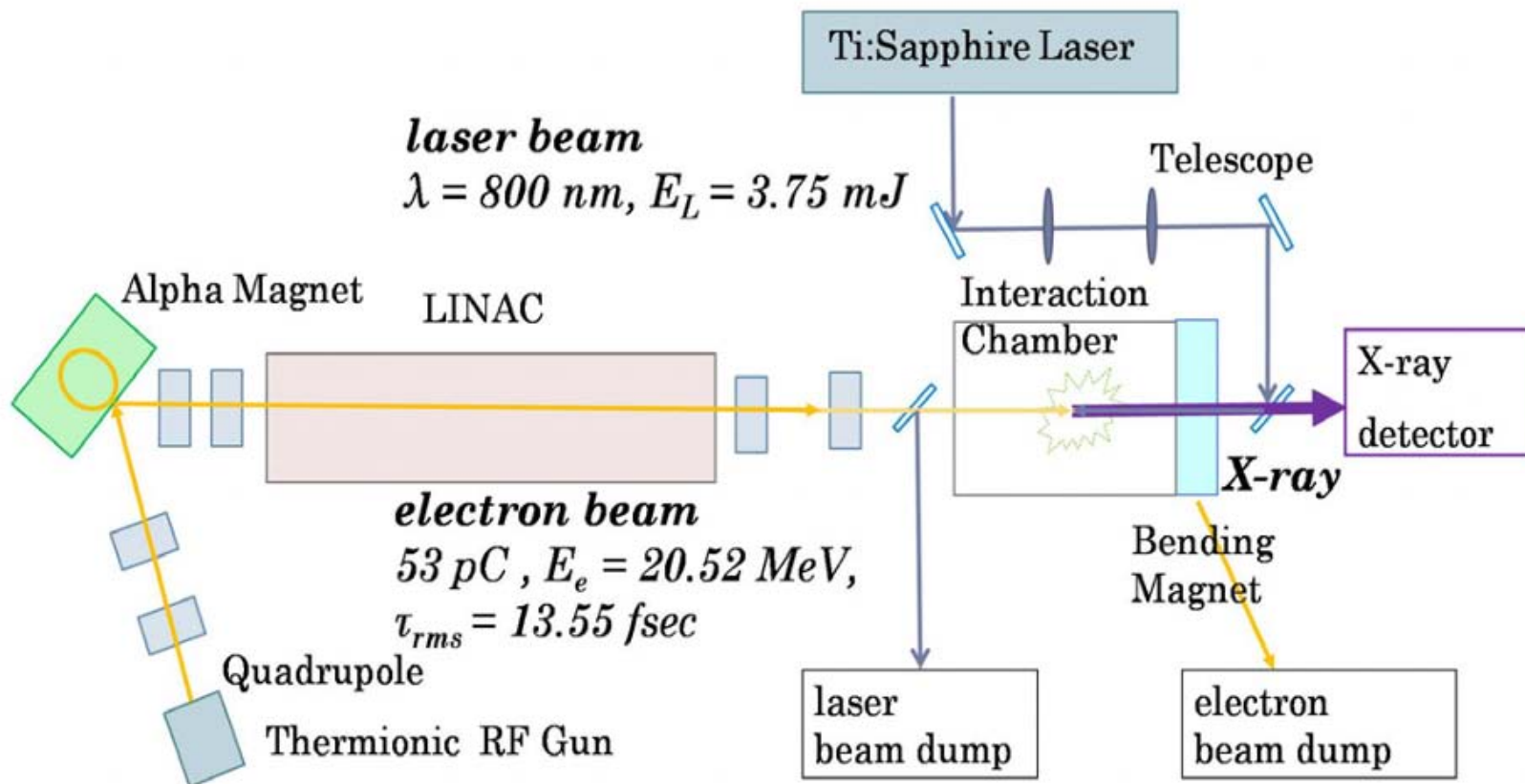
http://www.fhi-berlin.mpg.de/pc/PCres_methods.html

Examples:

- Ultrafast electron pulse radiolysis



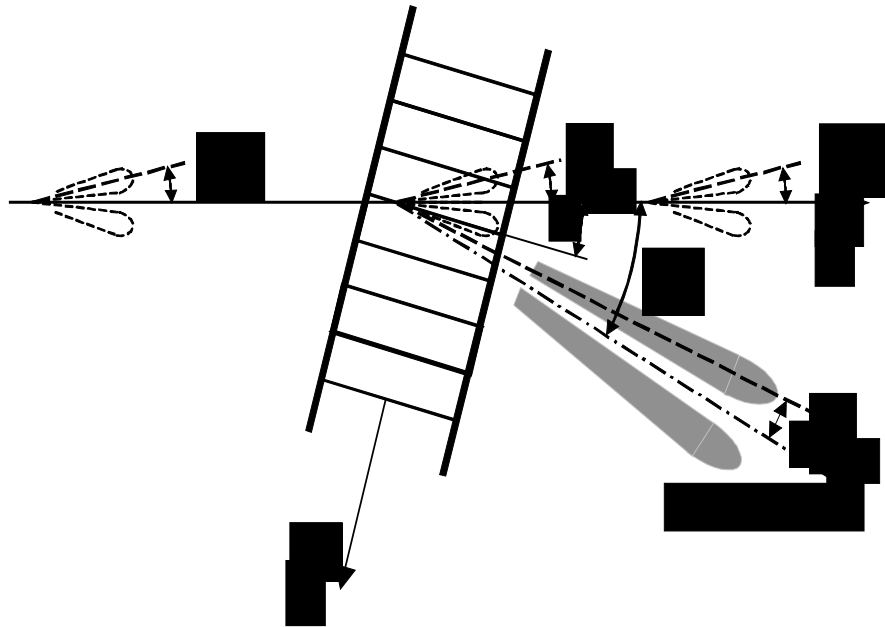
Examples: Head-on Inverse Compton scattering (ICS)

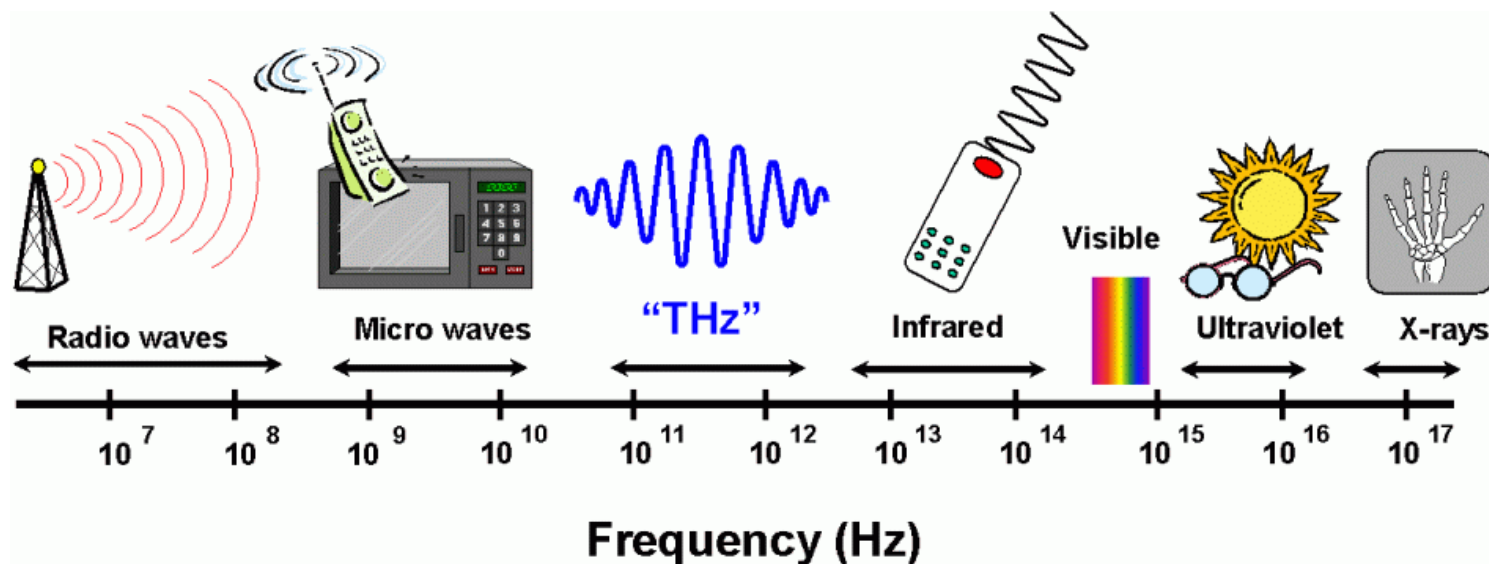


N.Y. Huang et al., Proceedings of PAC09, Vancouver, BC, Canada

Examples: Parametric X-rays

PXR as diffraction of virtual photons associated with relativistic charged particles

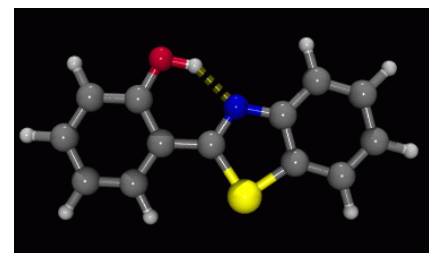
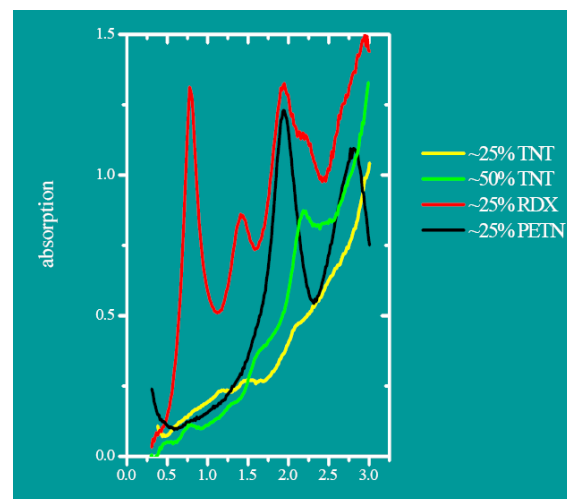




(<http://www.sp.phy.cam.ac.uk/SPWeb/research/thzcamera/WhatIsTHzImaging.htm>)

- gap between Microwaves and Infrared
- frequency from 100 GHz to 10 THz
- wave length from 30 mm – 0.3 cm
- In last 10-15 years, this area is unexplored region
 - lack of source and applications
- Now, THz technology is growing rapidly.

- Penetrate non-conducting material -
e.g. clothes, wood, plastic, ceramic, paper
- Blocked by metals, water
- Chemical sensitive
'finger print' absorption spectra
- Corresponds to intermolecular
vibration and rotation



Weak hydrogen bond

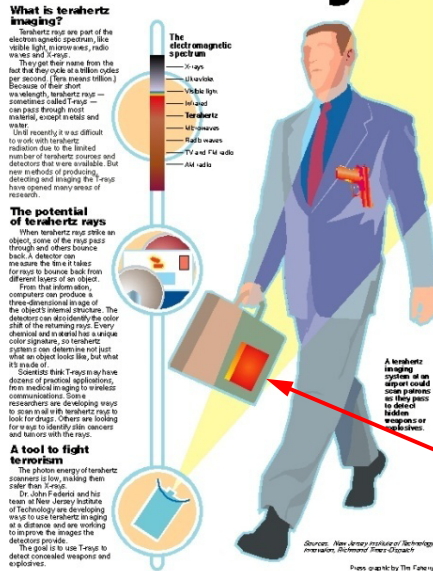
- ✦ **Safer:** low photon energies
- ✦ **Signature:** specific fingerprinting
- ✦ **More selective in soft materials**

e.g.,

- Security screening of hidden weapons or explosives
- Narcotics smuggling, suspicious packages
- Food quality control in frozen product
- Agriculture water-content check
- Medical diagnostic for cancerous cell
- Semiconductor wafer inspection

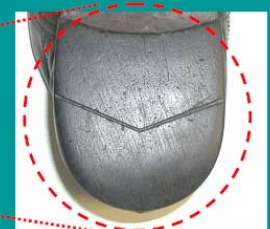
Defense applications

Tracking terror with T-rays



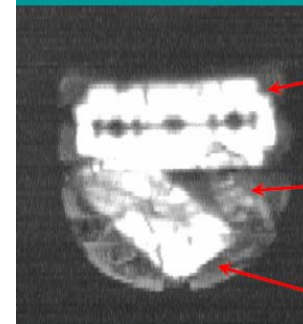
Standoff distance
detection of weapon
and explosives

Shoe Imaging



THz Image

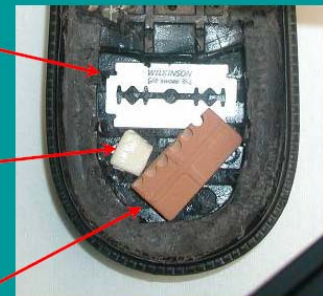
Visible Images



Razor blade

Plastic explosive

Ceramic



(<http://www.thznetwork.org/wordpress/index.php/thz-images>)

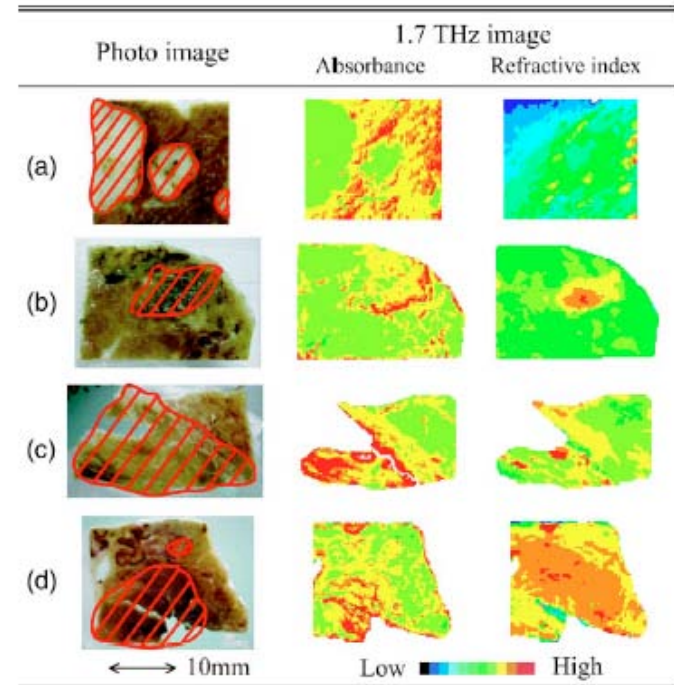
Medical applications

THz imaging diagnostic of cancer tissues

Nakagima et al. [Appl. Phy Let. 90 041102(2007)]



(<http://www.thznetwork.org/wordpress/index.php/thz-images>)



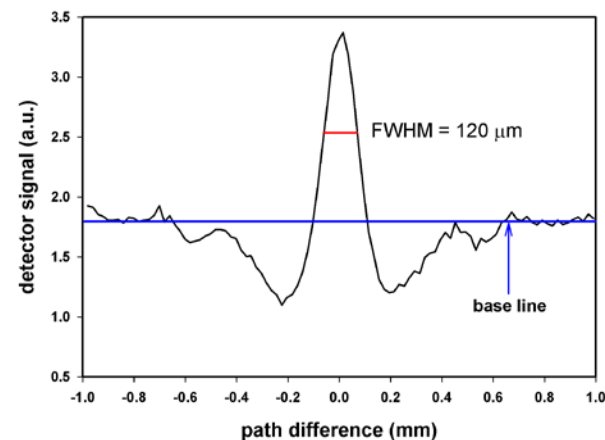
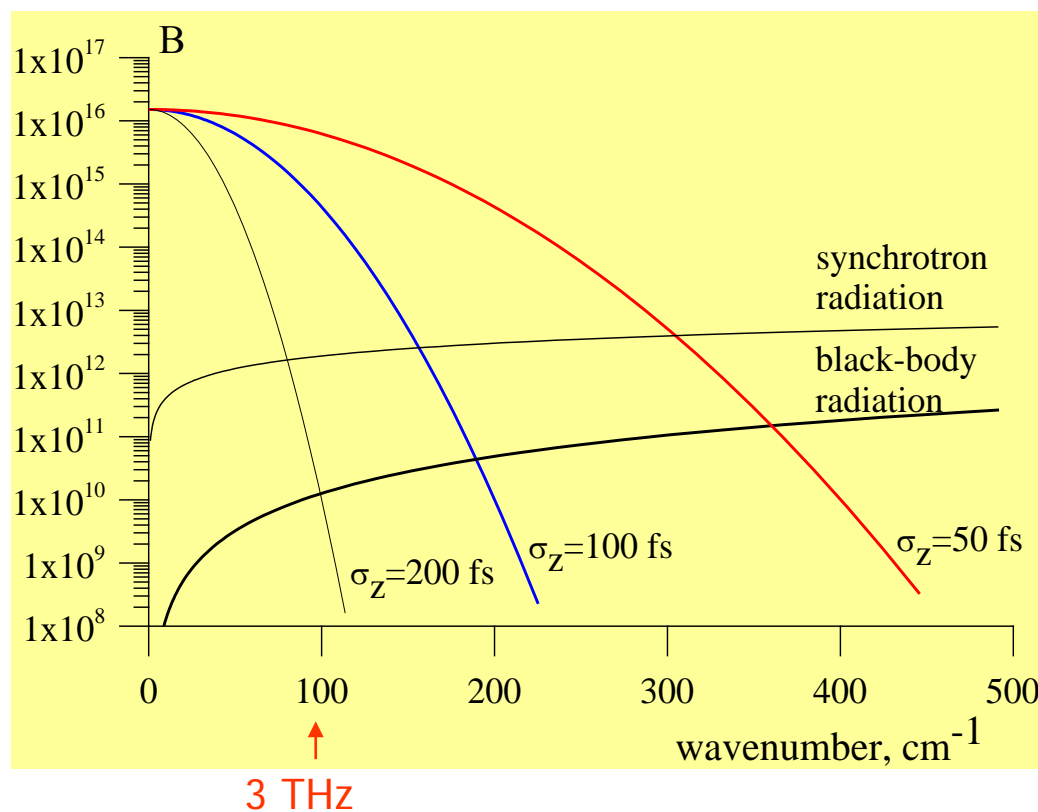
Inspection

Non-destructive THz imaging of drugs in an envelop

Kodo Kawase et al (RIKEN, Japan)

Fourier transform of the short electron bunch provides broad radiation spectrum

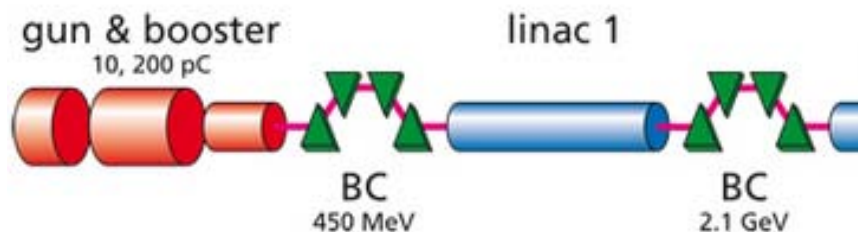
$$|\tilde{E}_0(\omega)|^2 = \tilde{E}_0(\omega)\tilde{E}_0^*(\omega)\alpha \int_{-\alpha}^{+\alpha} I_0(\delta)e^{i\omega\delta/c} d\delta = FT\{I_0(\delta)\}$$



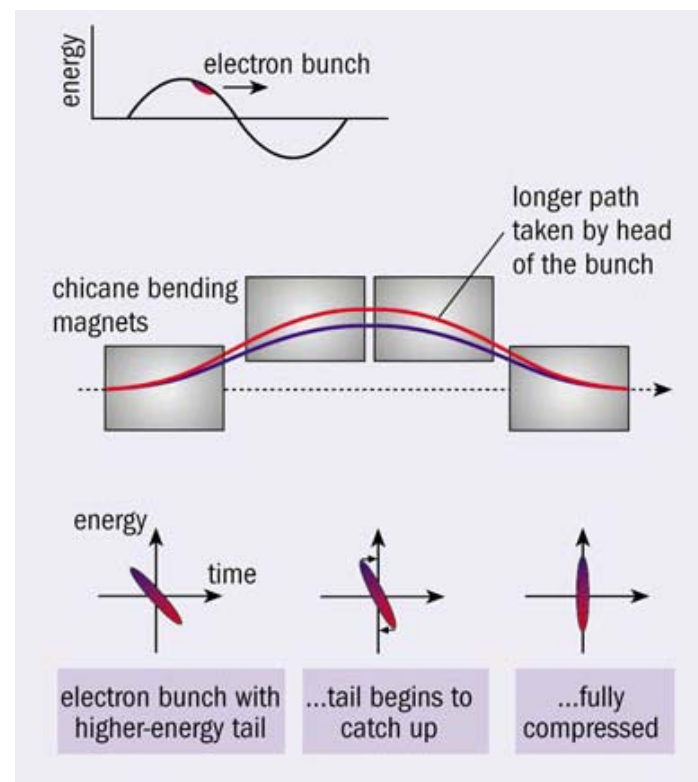
Calculated radiation brightness $B(\text{ph/s/mm}^2/100\% \text{BW})$ vs. wave number for CTR, SR and black body radiation.

Short electron bunches in the order of femtosecond scale can be generated by using an **RF gun** and a **magnetic bunch compressor**.

- High energy electron bunches with inversely linear correlation between energy-time distribution are normally compressed with **chicane**.

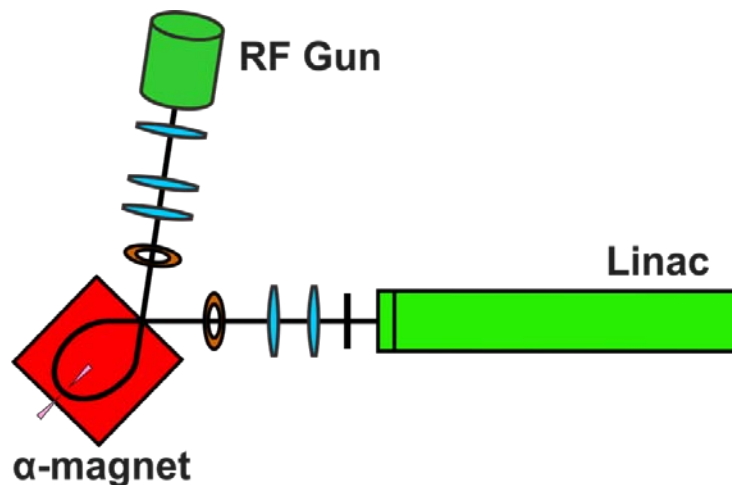


[http:// www.psi.ch/swissfel/swissfel-accelerator](http://www.psi.ch/swissfel/swissfel-accelerator)



<http://cerncourier.com/cws/article/cern/28925>

- Low energy electron bunches with linear correlation between energy-time distribution can be compressed with **α -magnet**.

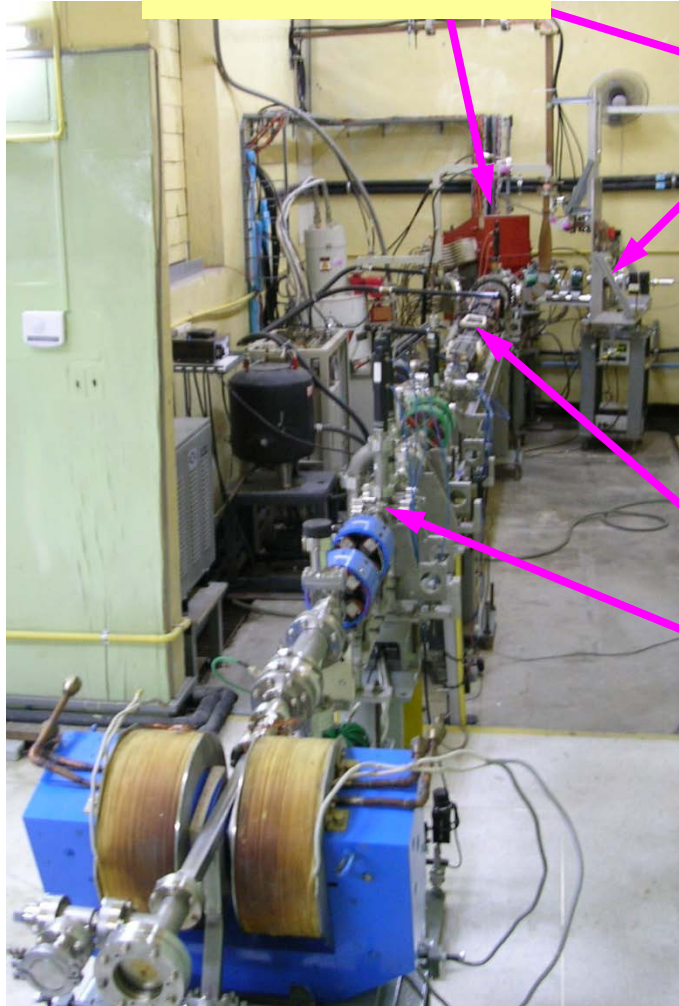


- Electrons with different momenta have different path lengths inside the α -magnet field.
- **Energy slits** are placed in a magnetic vacuum chamber to select only the core of the beam.
 - Emittance and energy spread are reduced.
 - The peak current is increased.

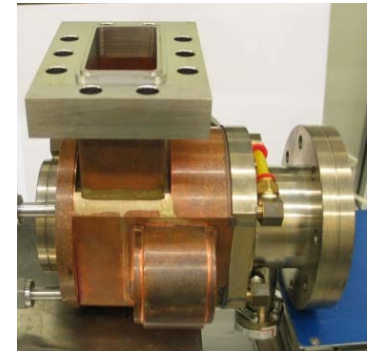
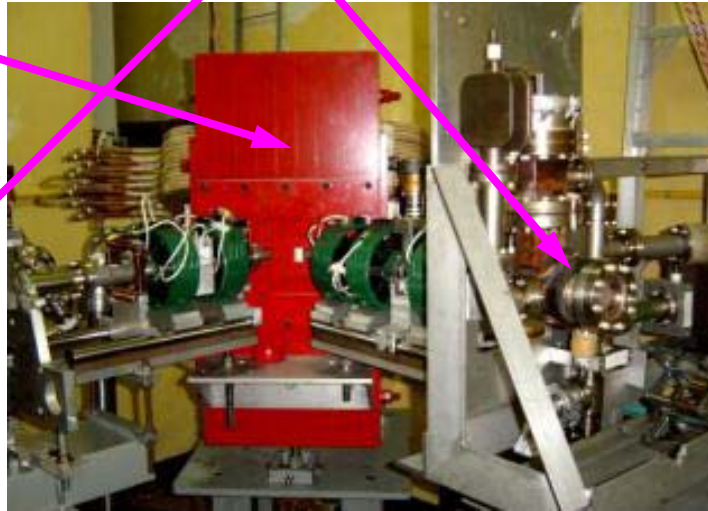




Alpha magnet
(bunch compressor)

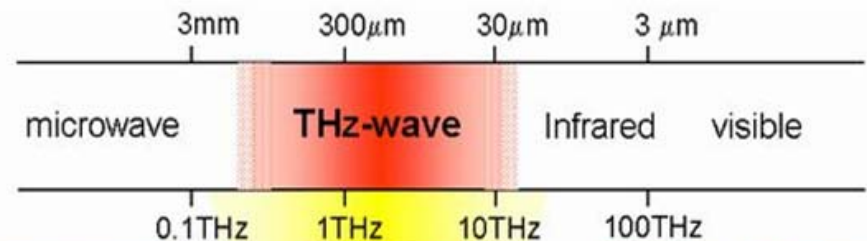


S-band thermionic RF-gun (2-2.5 MeV (30 MV/m @ cathode))



Linac (10-15 MeV (up to 30 MeV with higher RF power))

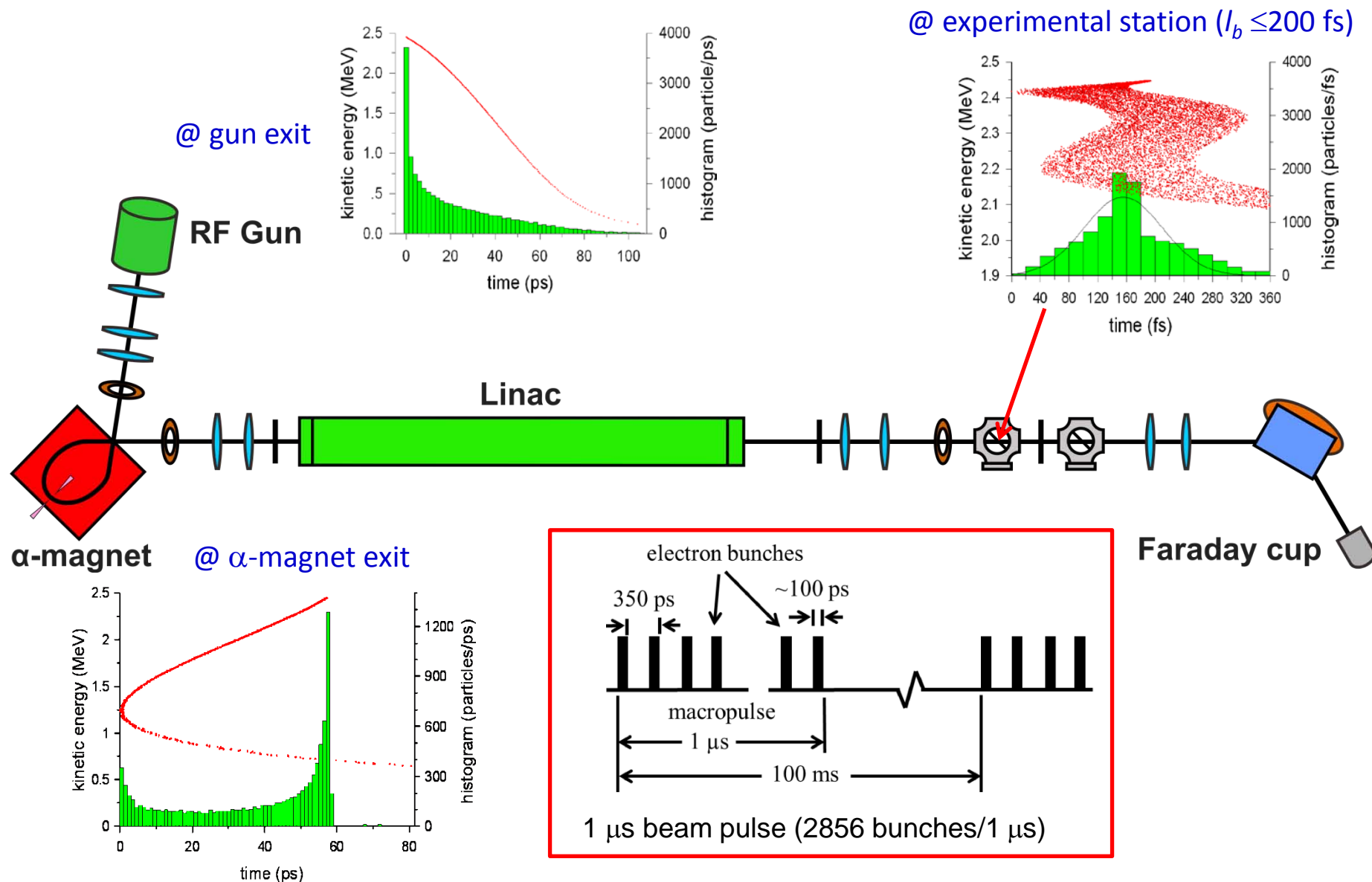
Transition radiation stations (0.23 nC/bunch, ≤ 200 fs)



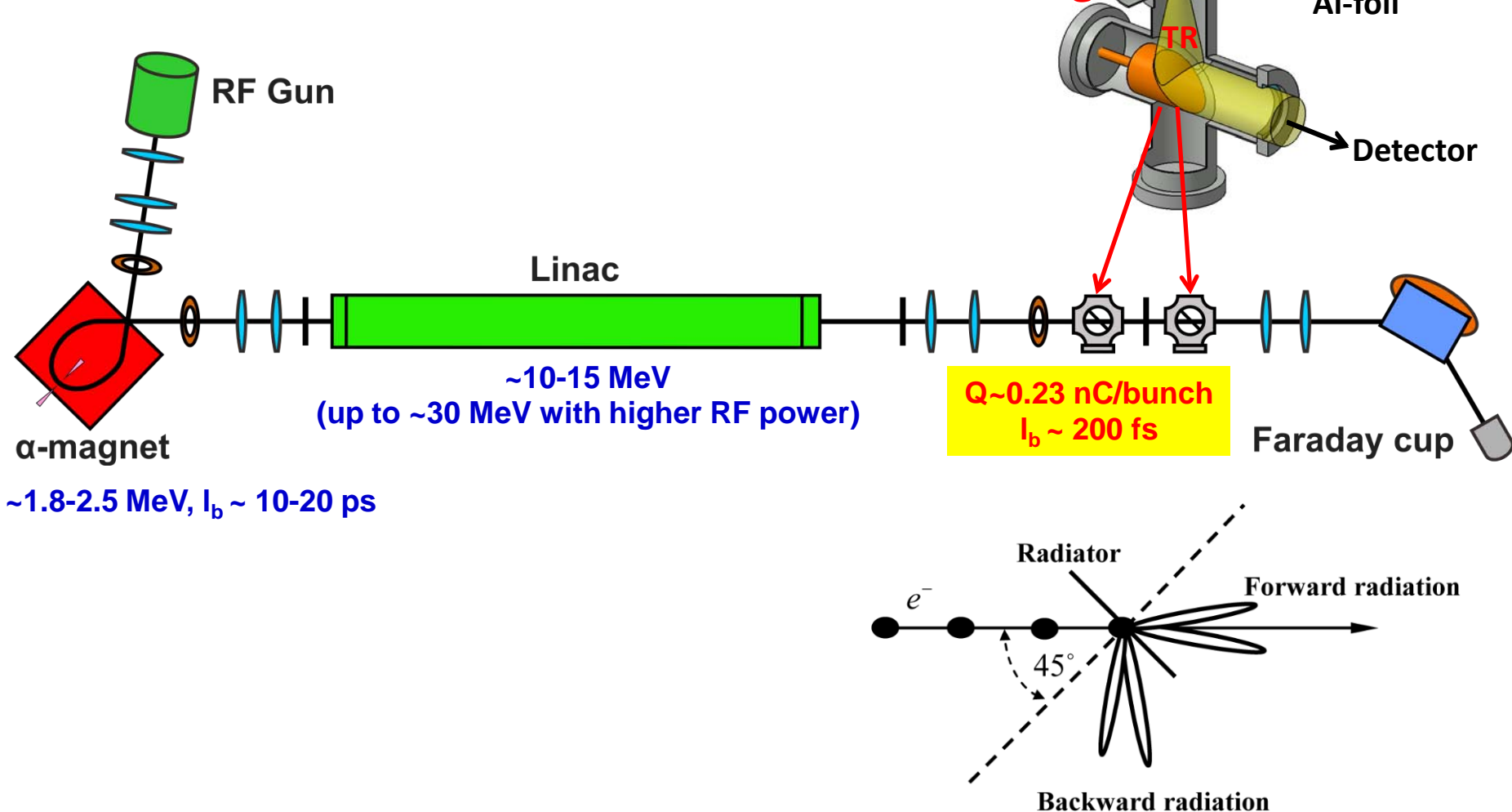
electronic approach

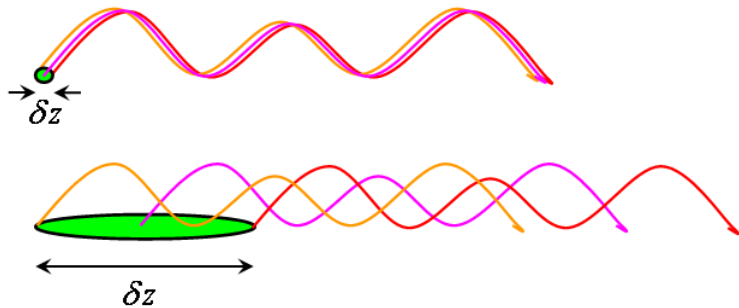
THz regime

photonic approach



$E_{\text{kin}} \sim 2\text{-}2.5 \text{ MeV}$ ($E_z \sim 30 \text{ MV/m}$ @ cathode)

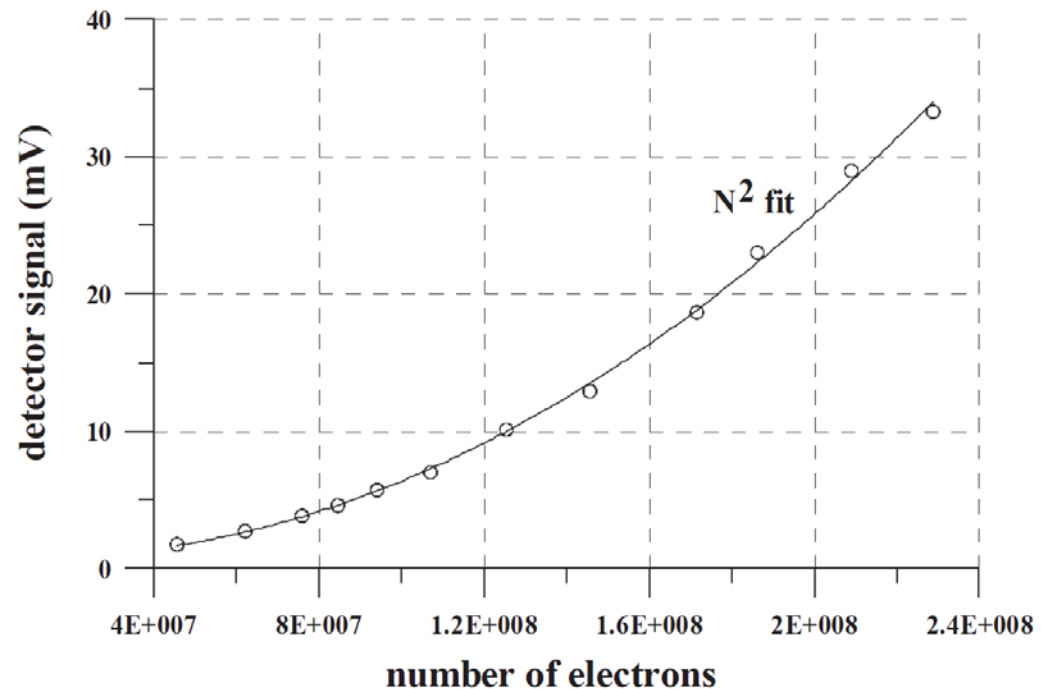


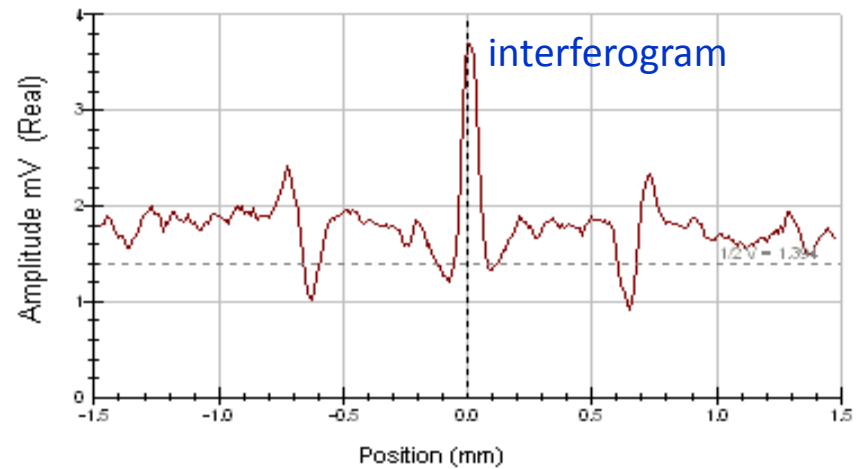
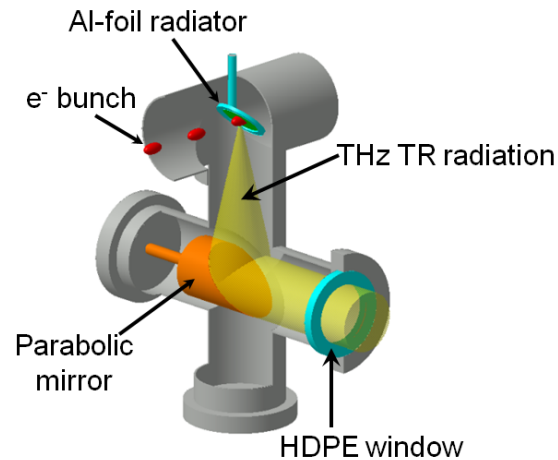


At wavelengths about or longer than the bunch length
↓
Radiation field add up coherently
↓
Electron short bunches is desired to produce coherent radiation
↓
Radiation intensity $\propto N^2$

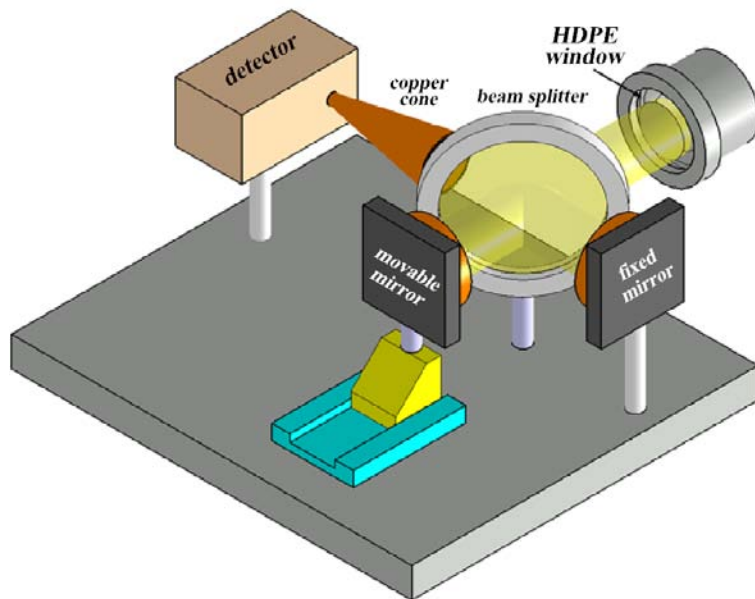
$$\frac{d^2 I}{d\omega d\Omega} = [N[1 - f(\omega)]^{ISR} + N^2 f(\omega)]^{CSR}$$

Detector signal of coherent transition radiation intensity vs. number of electrons.

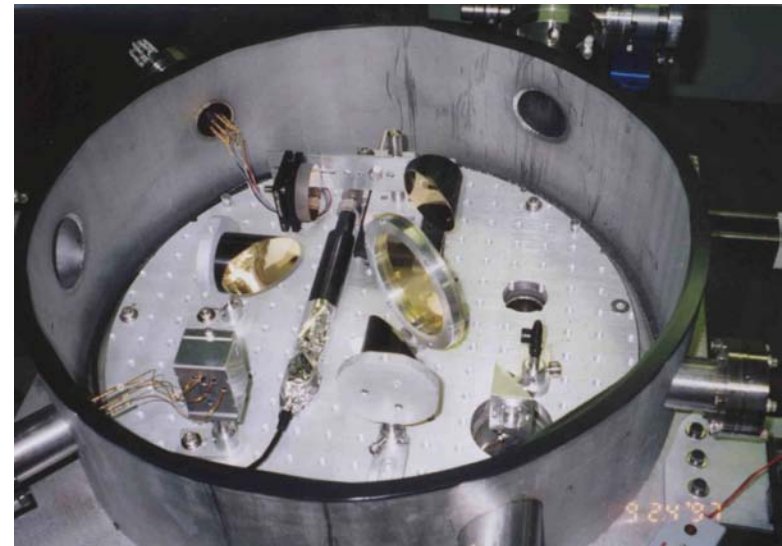




Measured electron bunch ~ 200 fs ($120 \mu\text{m}$)

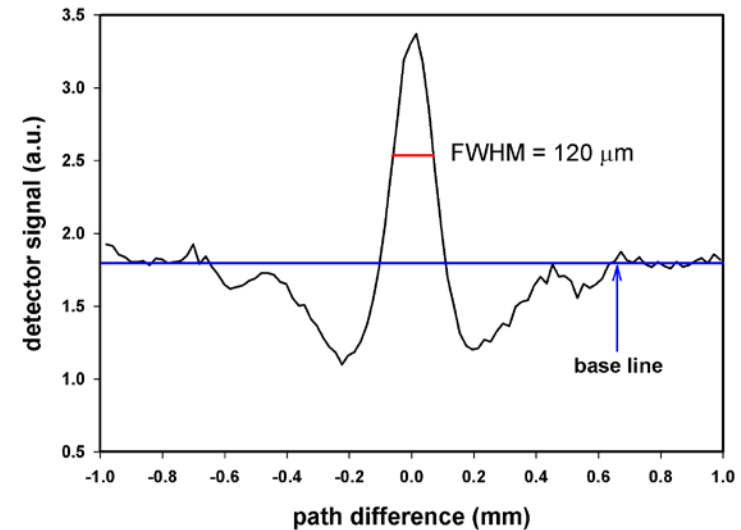
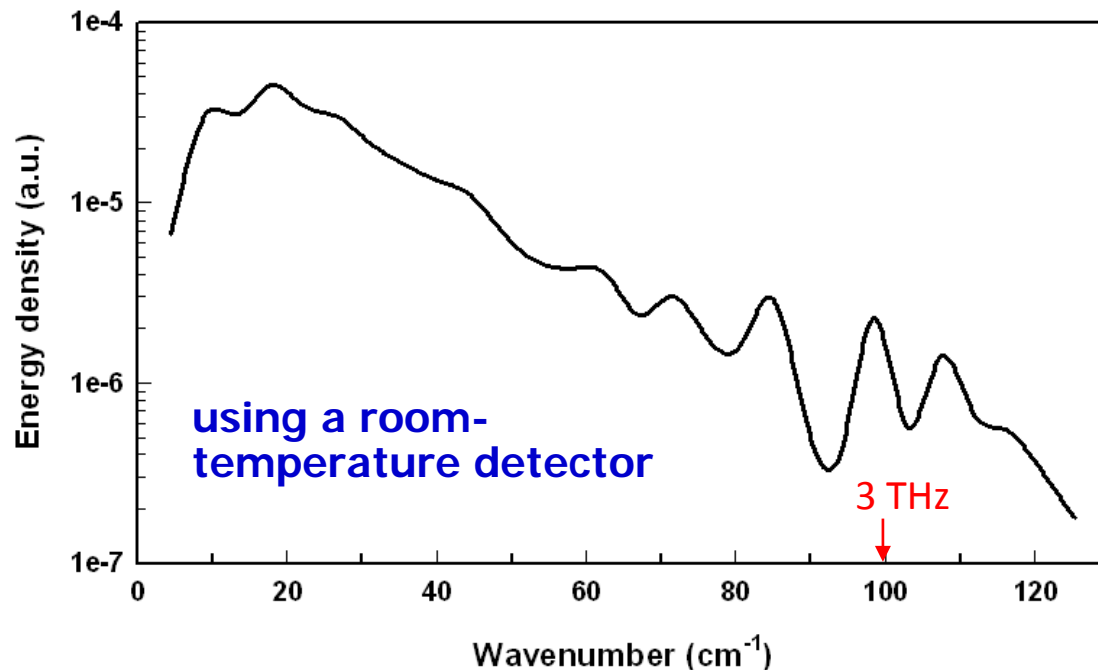


Michelson Interferometer



Fourier transform of the short bunch
provides broad radiation spectrum

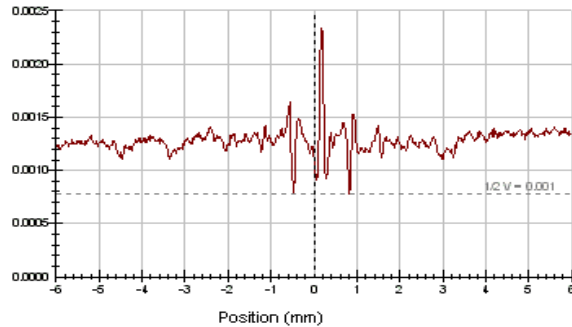
$$|\tilde{E}_0(\omega)|^2 = \tilde{E}_0(\omega)\tilde{E}_0^*(\omega)\alpha \int_{-\alpha}^{+\alpha} I_0(\delta)e^{i\omega\delta/c}d\delta = FT\{I_0(\delta)\}$$



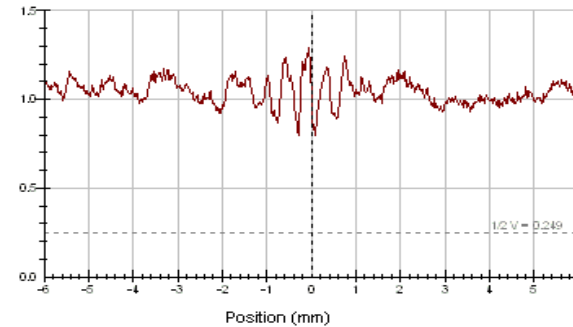
Measured electron bunch
~ 200 fs (120 μm)

Measured THz radiation covers
wave number of 5-80 cm⁻¹
(frequency range of 0.3 – 2.4 THz)

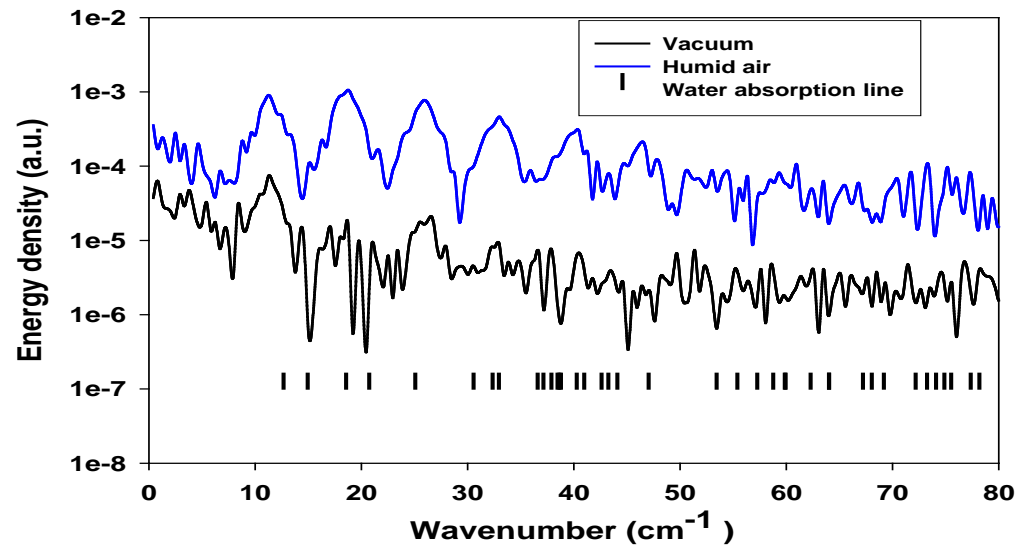
Water vapor spectrum & absorption lines

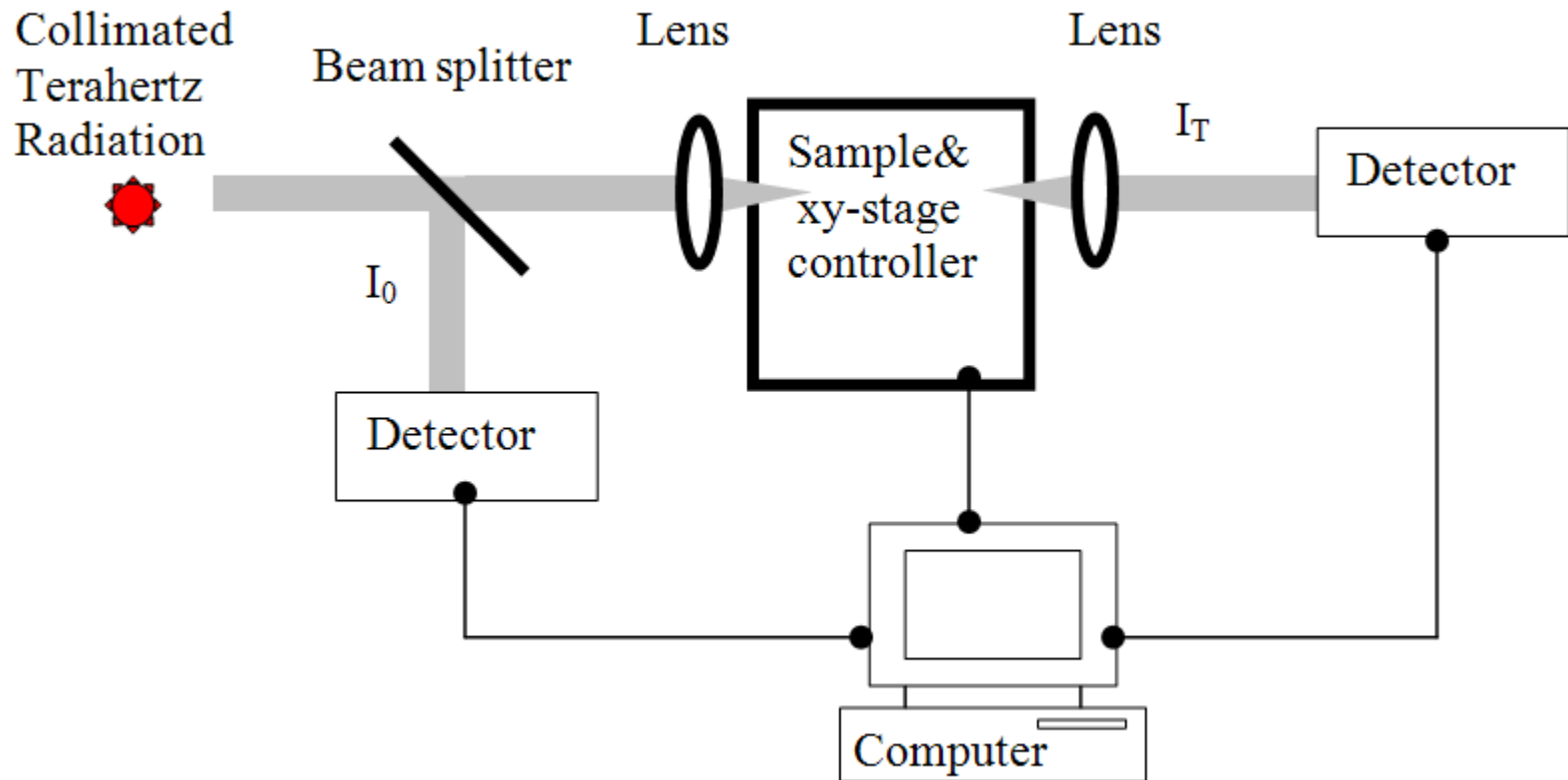


interferogram in vacuum

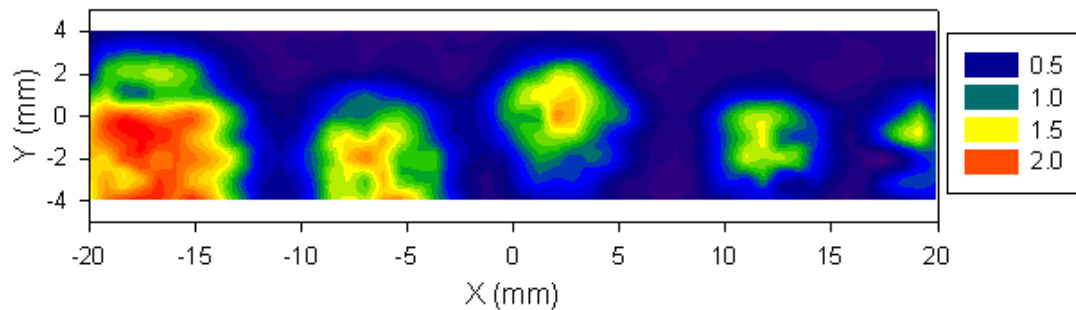
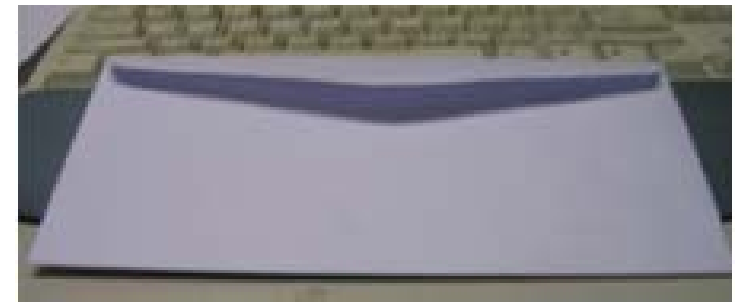
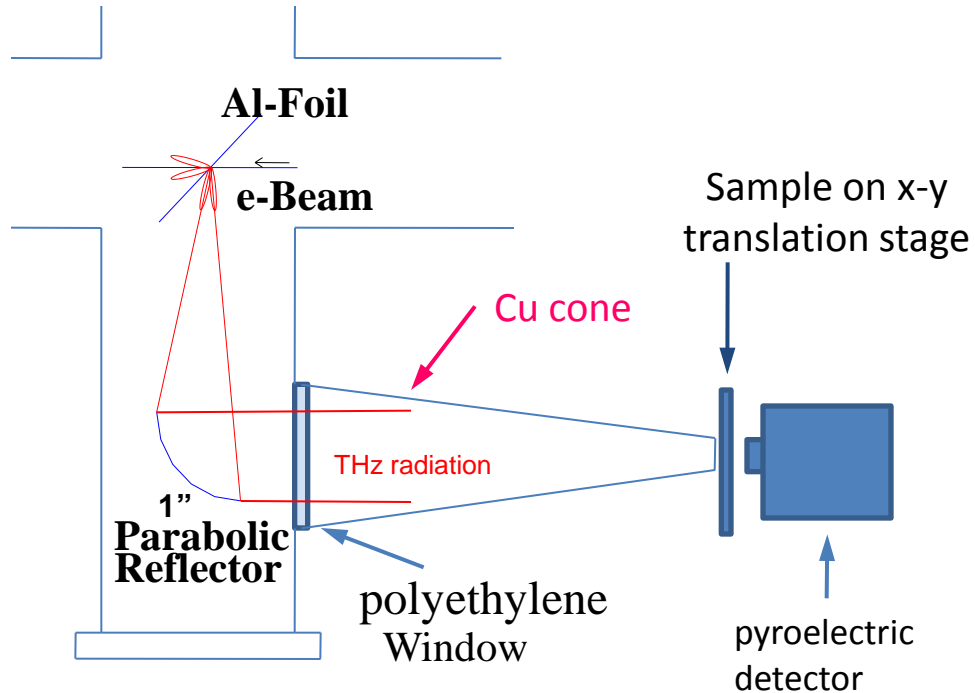


interferogram in humid air



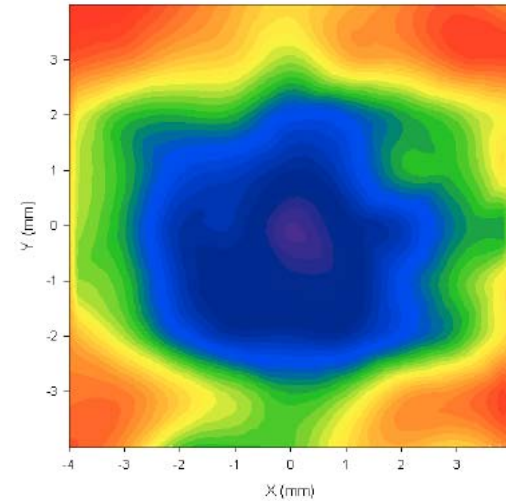
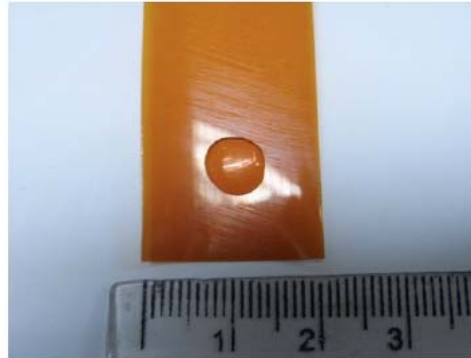


Design of THz imaging system at PBP, Chiang Mai University

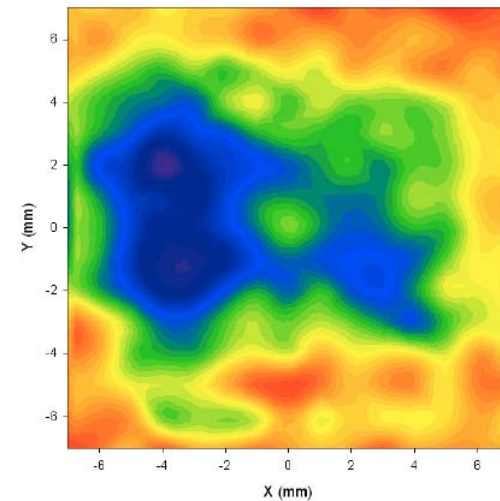


THz image of holes on Al-foil in an envelope

Water drop

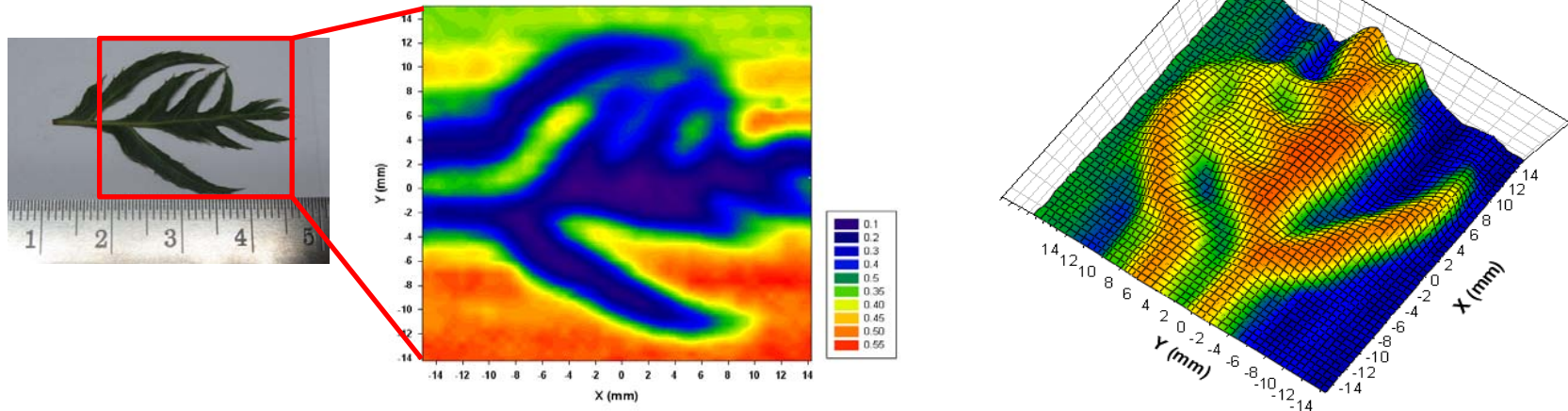


Cooked vs. Un-cooked Rice



THz image of live leaves (better resolution)

with 180x180 μm copper mesh filter (20–40 cm^{-1} transmission band)



Parameters	Value	Unit
Spectral range	5 - 80	cm^{-1} ,
Energy per macropulse (21.5% collection efficiency)	8.8	μJ
Polarization	radial	
Macropulse power	11	W
Microbunch power	19	kW
Average power (at 10 Hz rep. rate)	88	μW
Micropulse duration (σ_z)	200	fs
Micropulse separation	350	ps
Macropulse duration	0.8	μs
Number of radiation pulses/macropulse	2300	pulse

Electron beam parameters	
Resonant frequency	2856 MHz
Beam energy	10 MeV
Electron bunch length	200 fs
Bunch charge	230 pC
Macropulse length	1 μs
Repetition rate	10 Hz

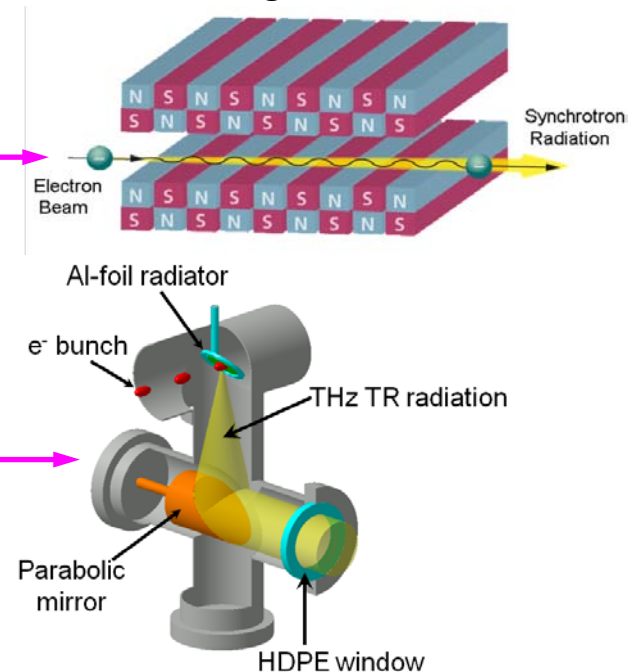
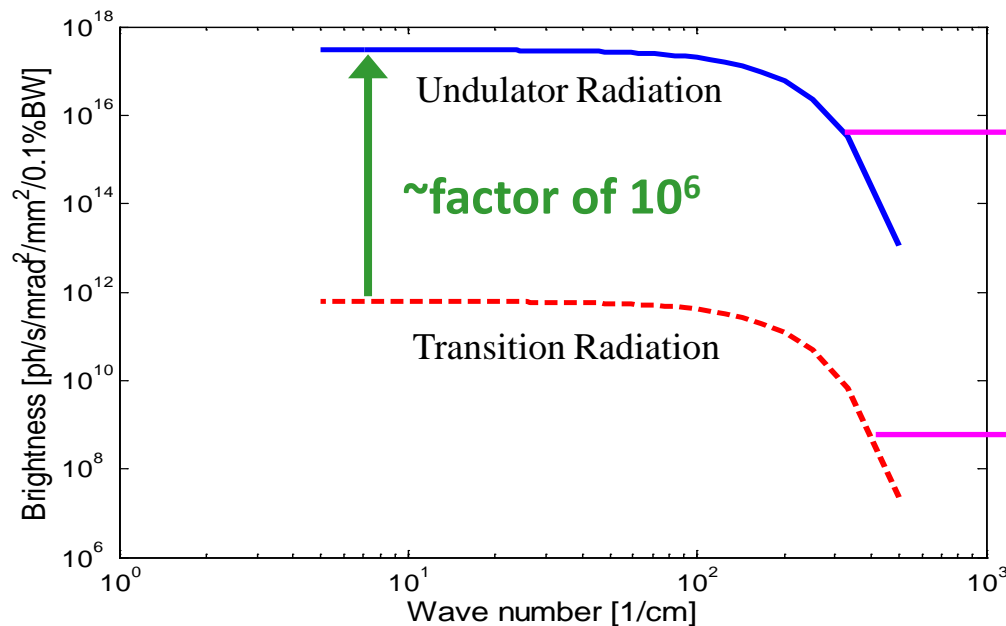
How to obtain
brighter radiation?

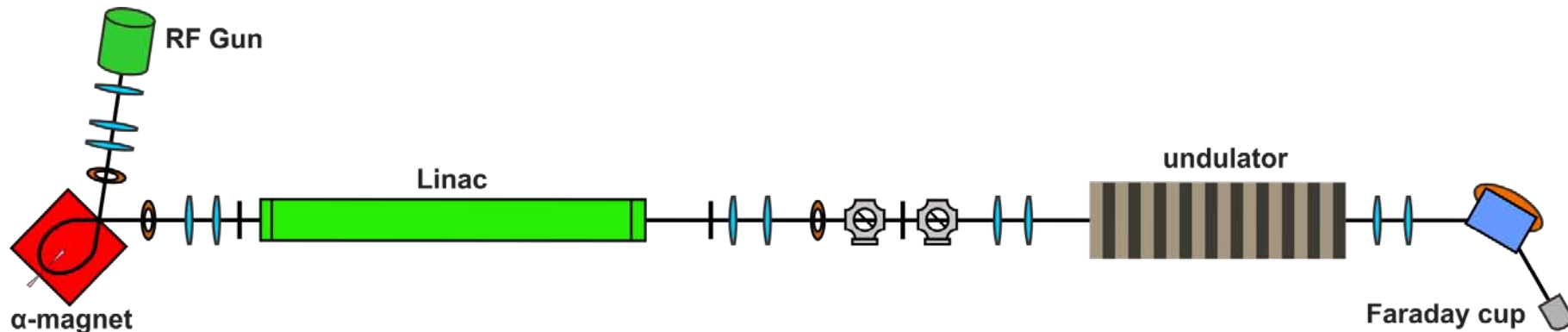
The greater the brightness, the more photons that can be concentrated on a spot.

Improvement of electron beam

Resonant frequency	2856 MHz
Beam energy	10 MeV
Electron bunch length	100 fs
Bunch charge	~100-200 pC
Macropulse length	1 μ s
Repetition rate	10 Hz

Brightness: number of photons produced per second, angular divergence of photons, cross-sectional area of beam, photons falling within a bandwidth (BW) of 0.1% of central wavelength.



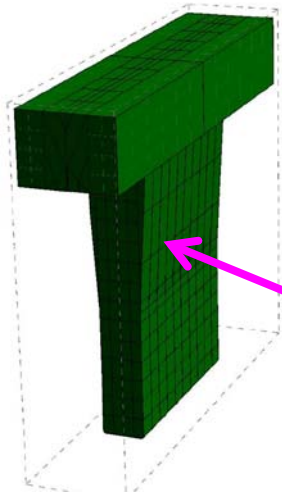


Planar Electromagnet Undulator (EMU)

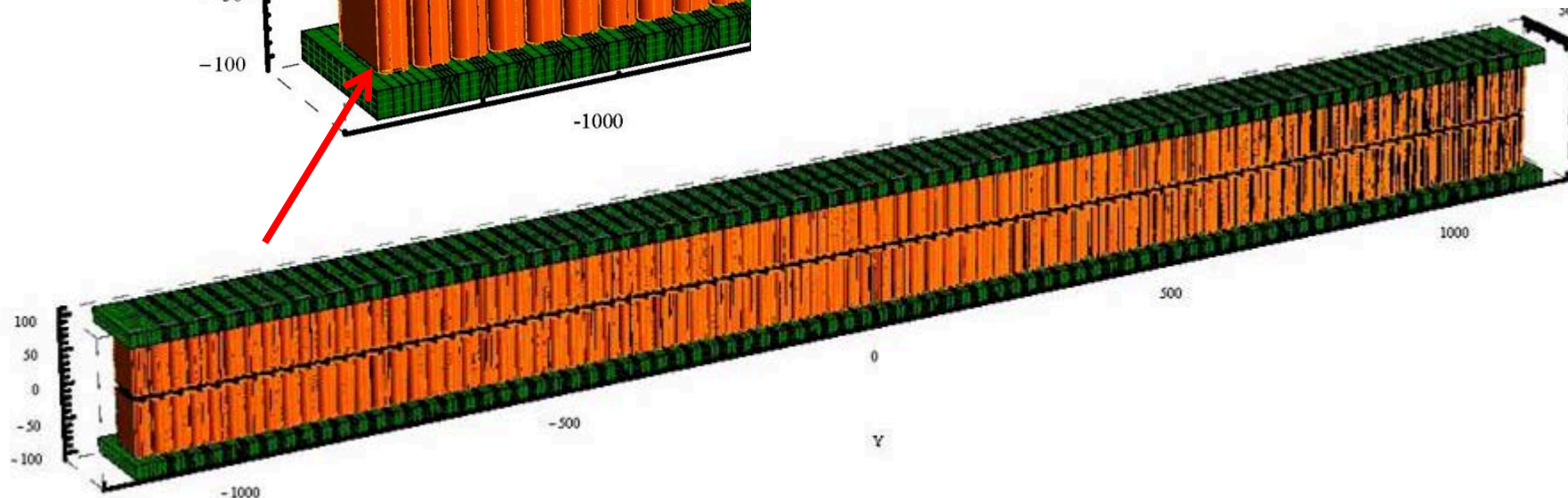
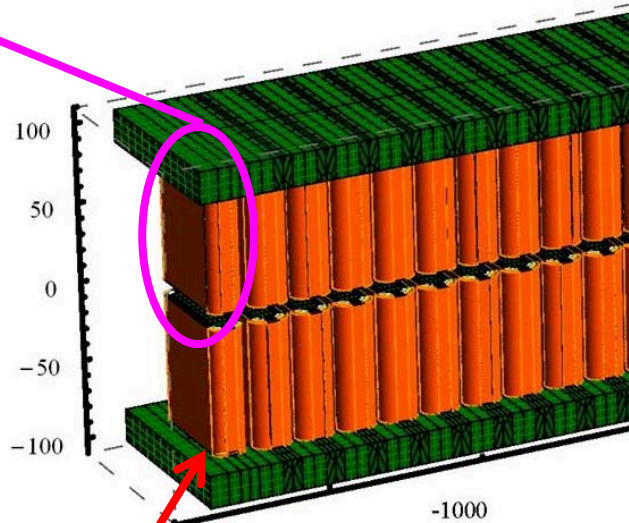
- Economical & simple choice
- Based on some experience of electromagnet dipoles and quadrupoles development in house
- Magnetic field strength is adjusted by varying the applied current
→ To avoid difficulty of mechanics for adjusting the gap

Preliminary parameters for CMU EMU

K-parameter	1- 1.2
Max. on-axis peak magnetic field	~0.2 T
Period length	55 mm
Undulator gap	10 mm
Number of periods	40
Radiation wavelength	~100 μm

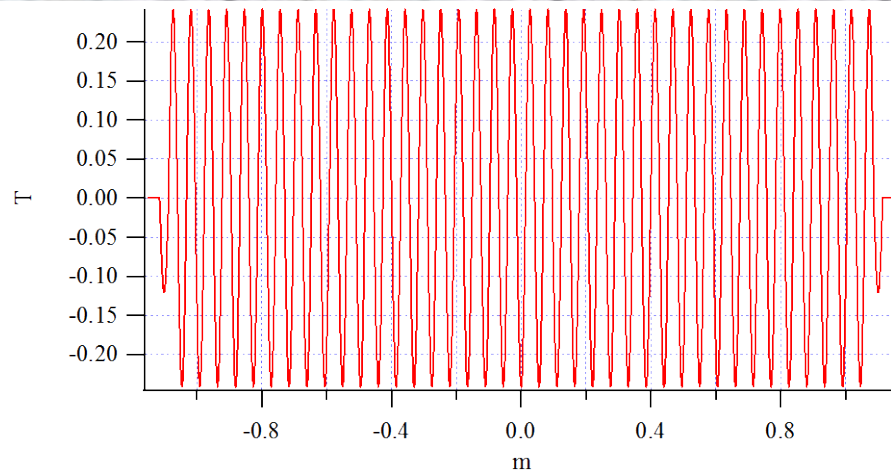


Pole optimization
for non-saturation
of magnetic field



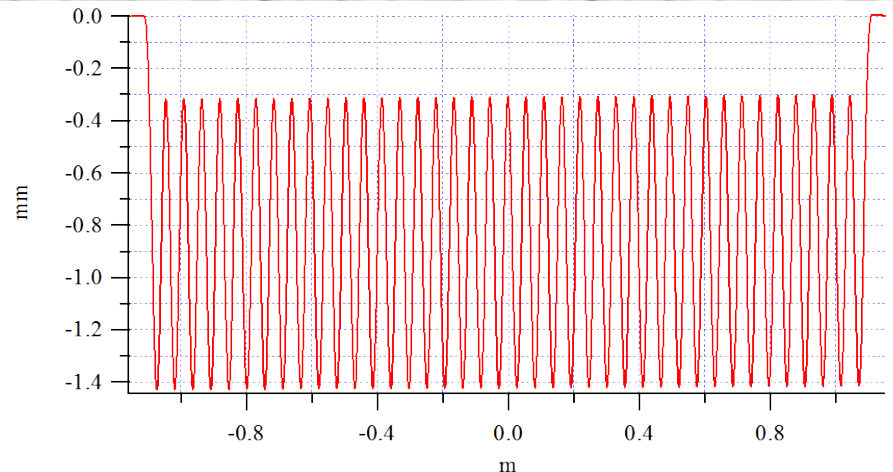
Design electromagnet undulator parameters

K-parameter	1.2
Max. on-axis peak magnetic field	0.23 T
Period length	55 mm
Undulator gap	10 mm
Number of periods	40
Radiation wavelength	~100 μm



B_y along distance in undulator (z)

$$B_0[T] = (2\pi mcK) / (e\lambda_u) = K / (0.934\lambda_u[cm])$$

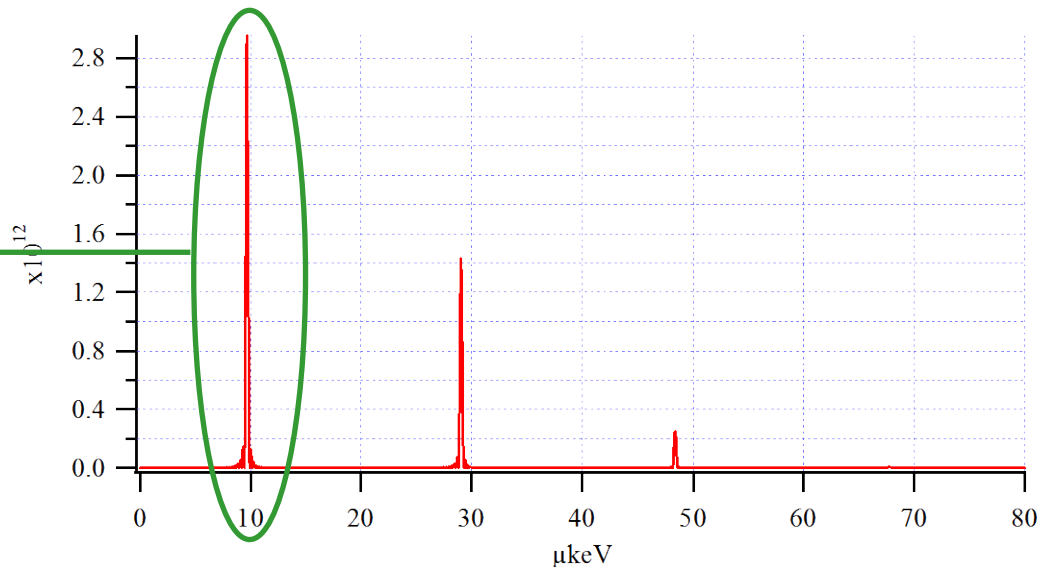


Transverse trajectory of electron along z

$$x = [K\lambda_u \sin(2\pi z / \lambda_u)] / [2\pi\beta\gamma]$$

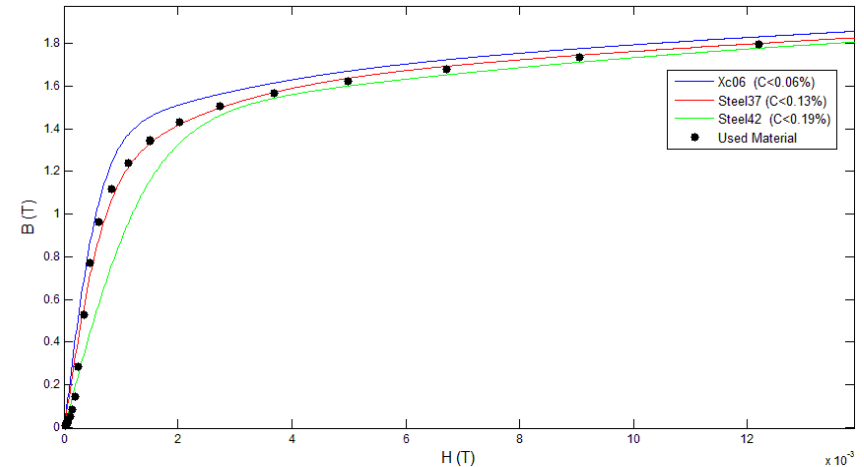
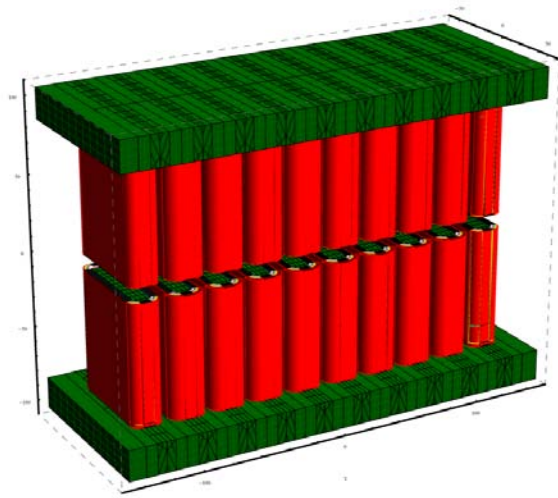
Fundamental frequency
~ 2.4 THz

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \theta^2\gamma^2 \right)$$

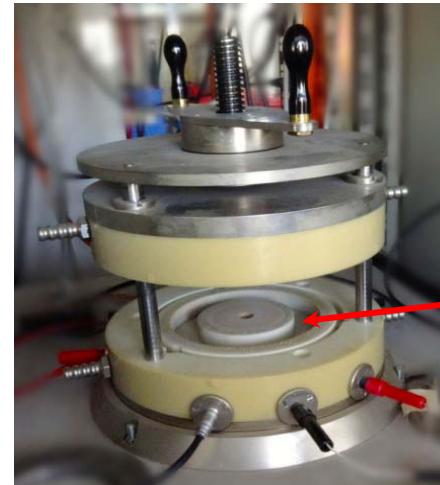


B2E simulated undulator spectrum

RADIA model (8 main poles + 2 end poles)



B-H curve of the material used to construct the undulator was measured by using **Split-coil permeameter** @ DESY



Characteristics of particle beams and reliability in operation of an accelerator facility strongly depend on properties of the source.

- Development of high-brightness electron beams is a key and a critical issue in the success of most electron accelerator projects.
- RF electron guns are the powerful source with specifications depending on types of accelerators and applications.
- RF gun technology is one of main streams in the development of particle accelerators.

**Thank you for your
attention!**

Questions & Discussion